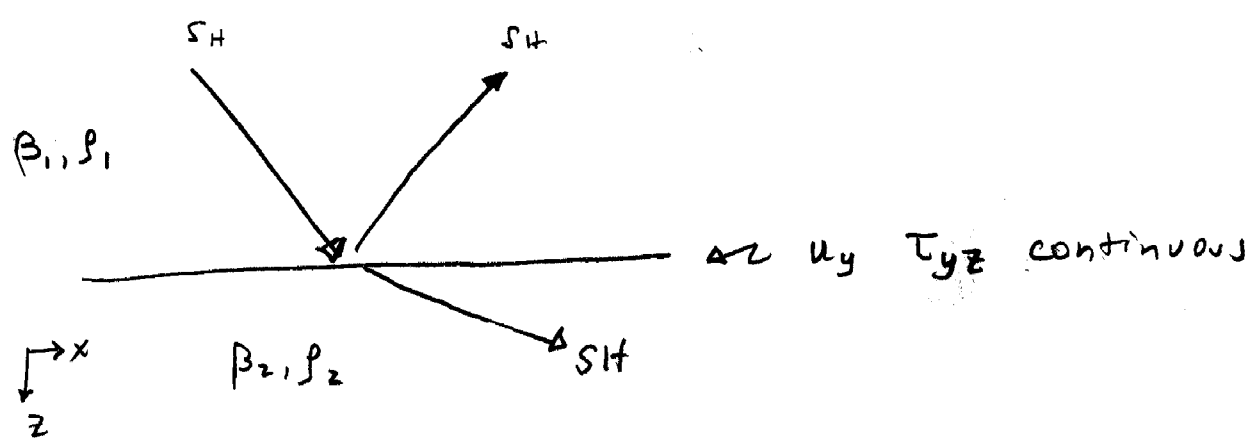


(1)



TOP $u_y^{(1)} = \check{C}_1 \cos(\omega p_x x + \omega p_z^{(1)} z - \omega t) + \check{C}'_1 \cos(\omega p_x x - \omega p_z^{(1)} z - \omega t)$

BOT $u_y^{(2)} = \check{C}_2 \cos(\omega p_x x + \omega p_z^{(2)} z - \omega t)$

p_x equal for all waves. p_z 's given by snell's law

$$p_z^{(1)} = \sqrt{\beta_1^{-2} - p_x^2} \quad p_z^{(2)} = \sqrt{\beta_2^{-2} - p_x^2}$$

since $\frac{\sin \theta_1}{\beta_1} = p_x^{(1)}$ and $\frac{\cos \theta_1}{\beta_1} = p_z^{(1)}$

square and add to get

$$\frac{1}{\beta_1^2} (\sin^2 \theta + \cos^2 \theta) = p_x^2 + p_z^{(1)2}$$

$$\tau_{ij} = \lambda \delta_{ij} \sum_p \epsilon_{pp} + 2\mu \epsilon_{ij} \quad \tau_{yz} = 2\mu \epsilon_{yz}$$

$$\epsilon_{ij} = \frac{1}{2} \left(\frac{du_i}{dx_j} + \frac{du_j}{dx_i} \right) \quad \frac{1}{2} \frac{du_y}{dz} = \epsilon_{yz}$$

and $\tau_{yz} = \mu \frac{du_y}{dz}$

TOP: $-\mu_1 \check{C}_1 \omega p_z^{(1)} \sin(\cdot) + \mu_1 \check{C}'_1 \omega p_z^{(1)} \sin(\cdot)$

BOT: $-\mu_2 \check{C}_2 \omega p_z^{(2)} \sin(\cdot)$

Continuity of displacement $u_y^{(1)}(z=0) = u_y^{(2)}(z=0)$

(2)

$$\dot{C}_1 \cos(\omega P_x x - \omega t) + \dot{C}'_1 \cos(\omega P_x x - \omega t) = \dot{C}_2 \cos(\omega P_x x - \omega t)$$

$$\dot{C}_2 - \dot{C}'_1 = \dot{C}_1$$

Continuity of traction $\tau_{yz}^{(1)}(z=0) = \tau_{yz}^{(2)}(z=0)$

$$-\mu_1 \dot{C}_1 \omega P_z^{(1)} \sin(\omega P_x x - \omega t) + \mu_1 \dot{C}'_1 \omega P_z^{(1)} \sin(\cdot) = -\mu_2 \dot{C}_2 \omega P_z^{(2)} \sin(\cdot)$$

$$+\mu_2 P_z^{(2)} \dot{C}_2 + \mu_1 P_z^{(1)} \dot{C}'_1 = +\mu_1 P_z^{(1)} \dot{C}_1$$

$$\begin{bmatrix} 1 & -1 \\ \mu_2 P_z^{(2)} & \mu_1 P_z^{(1)} \end{bmatrix} \begin{bmatrix} \dot{C}_2 \\ \dot{C}'_1 \end{bmatrix} = \begin{bmatrix} 1 \\ \mu_1 P_z^{(1)} \end{bmatrix} \dot{C}_1$$

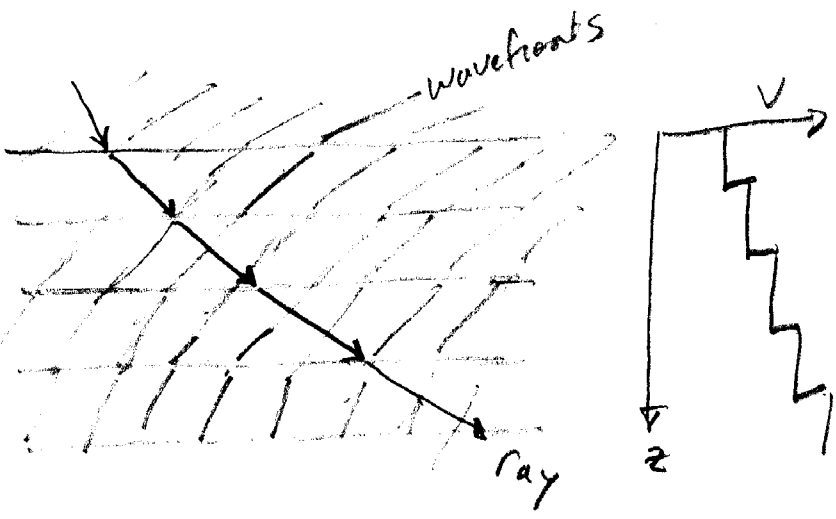
$$\det = \mu_1 P_z^{(1)} + \mu_2 P_z^{(2)}$$

$$\begin{bmatrix} \dot{C}_2 \\ \dot{C}'_1 \end{bmatrix} = \frac{1}{\det} \begin{bmatrix} \mu_1 P_z^{(1)} & 1 \\ -\mu_2 P_z^{(2)} & 1 \end{bmatrix} \begin{bmatrix} 1 \\ \mu_1 P_z^{(1)} \end{bmatrix} \dot{C}_1$$

$$= \begin{bmatrix} 2\mu_1 P_z^{(1)} / [\mu_1 P_z^{(1)} + \mu_2 P_z^{(2)}] \\ [\mu_1 P_z^{(1)} - \mu_2 P_z^{(2)}] / [\mu_1 P_z^{(1)} + \mu_2 P_z^{(2)}] \end{bmatrix} \dot{C}_1$$

③

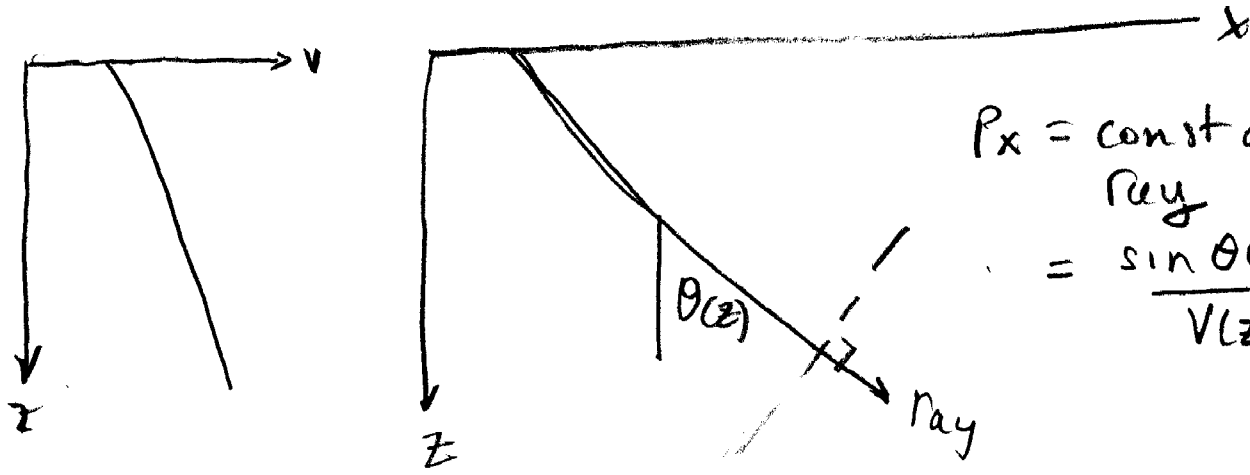
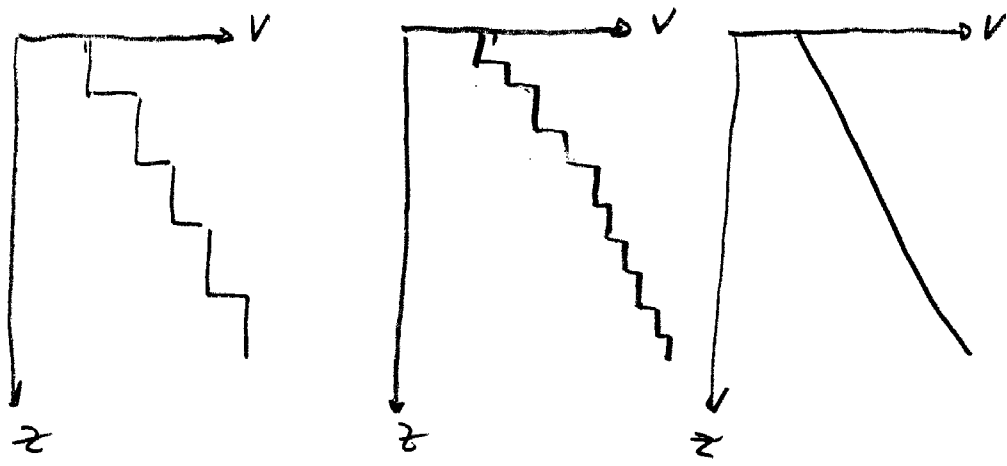
Vertically-stratified media.



P_x equal
in all layers
just draw n 's
connect to make
"ray"

P_x constant along The "ray".

Limit of lots of fine layers

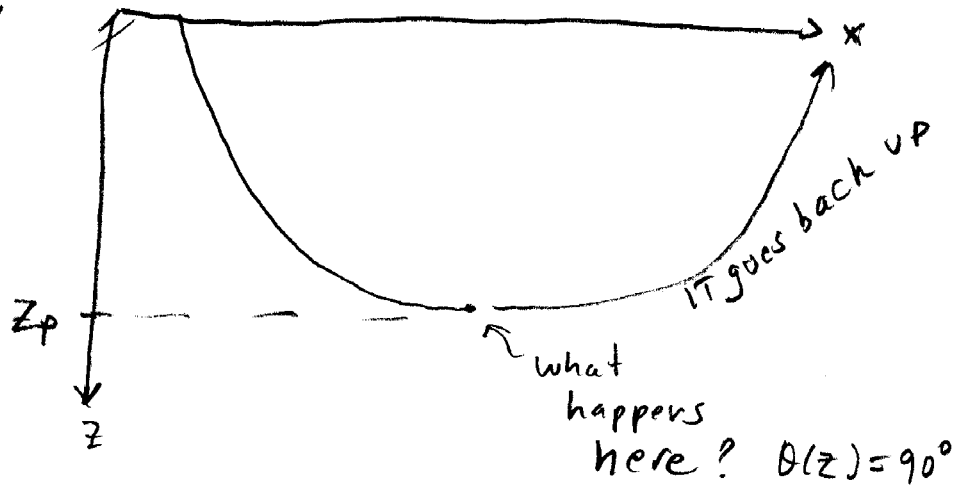
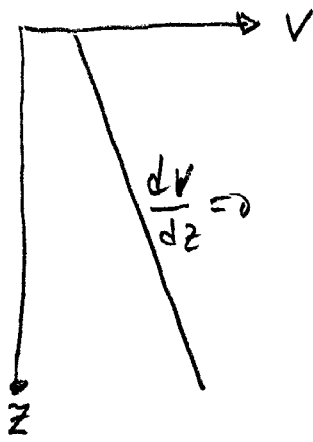


$$P_x = \text{const along Ray} = \frac{\sin \theta(z)}{V(z)}$$

wavefronts normal to ray

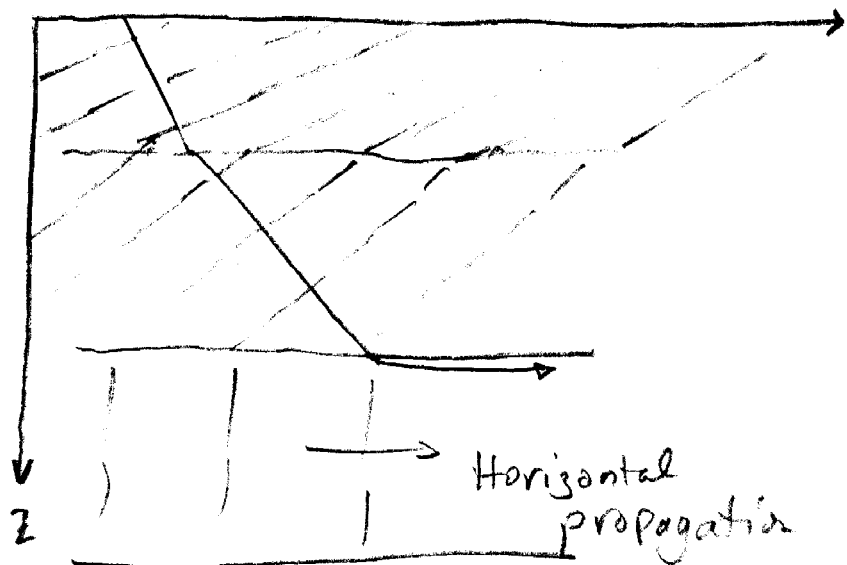
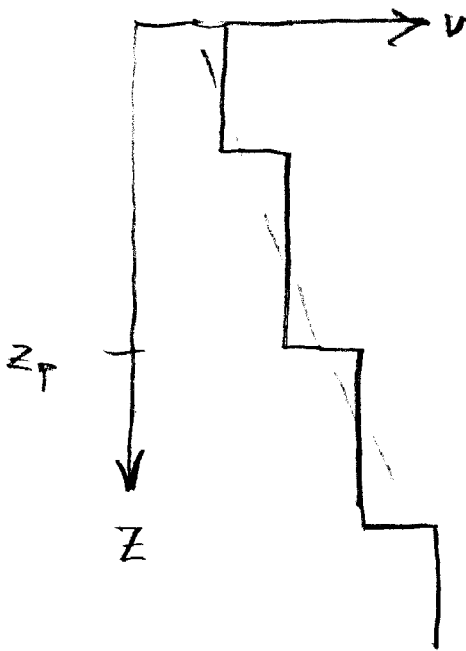
4

Case: velocity increase with depth, rays curve up



$$P_x = \frac{\sin(\theta(z))}{v(z)}$$

$P_x = v'(z_p)$ turning point of ray

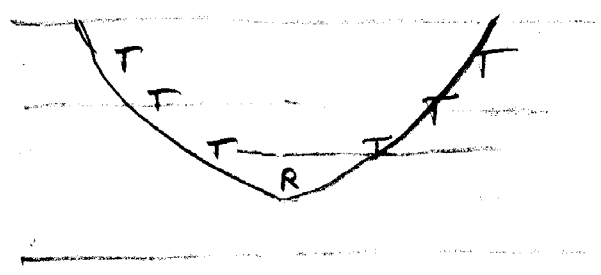
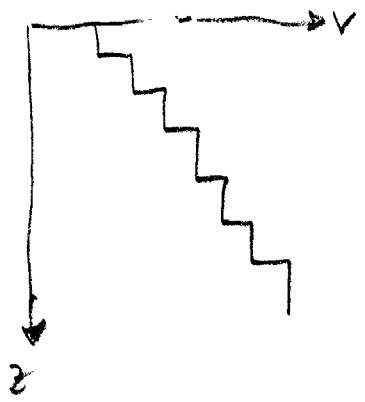


? what happens here

$$P_z = \sqrt{1/v^2 - P_x^2} \text{ is imaginary}$$

$$\cos(\omega P_z z) = \cosh(i\omega P_z z)$$

← exponential decay



most important waves
 Transmitted, except for
 lowest layer in which
 case wave field is
 reflected.

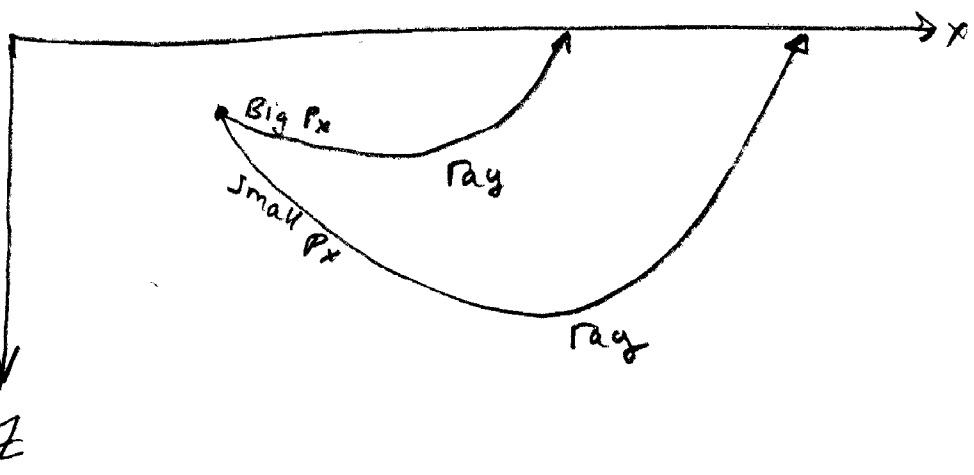
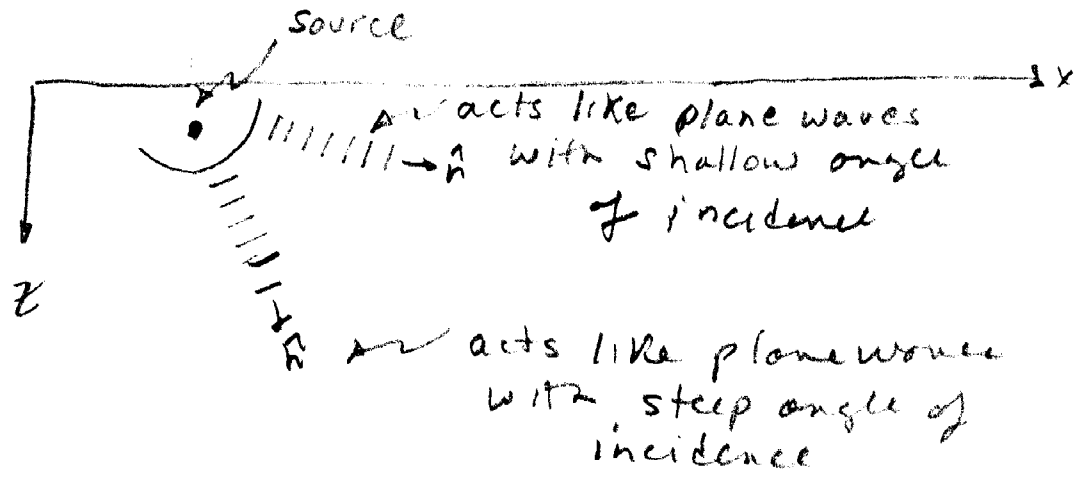
Ray Theory : smoothly varying media
 not necessarily vertically stratified
 λ short compared to
 scale-length of velocity
 variability.

Guiding Principle: "Almost planar waves"

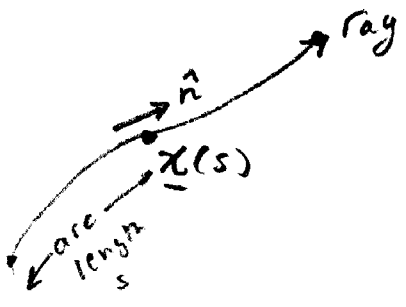
locally, The wave front moves
 parallel to itself at a speed
 given by the local velocity, v

$T(\underline{x}_1, \underline{x}_2)$ traveltime from \underline{x}_2 to \underline{x}_1 .

$$\nabla_{\underline{x}} T = \frac{\hat{n}}{v}$$



The equation for the ray is completely specified by the condition $\nabla T = \hat{n}/v$



proof

start with $\frac{\hat{n}}{v} = \nabla T$

dot with self $\frac{1}{v^2} = \nabla T \cdot \nabla T$

take $\frac{d}{ds}$ of $\frac{\hat{n}}{v} = \nabla T$

note $\hat{n} = \frac{dx}{ds}$

$$\frac{d}{ds} (v^{-1} \hat{n}) = \frac{d}{ds} \left(v^{-1} \frac{dx}{ds} \right) = \frac{d}{ds} \nabla T = \hat{n} \cdot \nabla (\nabla T)$$

$$= \frac{dx}{ds} \cdot \nabla (\nabla T) = v \nabla T \cdot \nabla \nabla T = \frac{v}{v^2} \nabla (\nabla T^2)$$