# Solid Earth Dynamics 

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Lecture 24

## Glaciology



How fast does ice melt?

$$
1360 \frac{\mathrm{~J}}{\mathrm{~m}^{2} \mathrm{~s}}
$$

albedo $\alpha=0.9$ but less that far north, say $1000 \frac{\mathrm{~J}}{m^{2} S}$
glacier at 0 degC
Heat of fusion of ice $3.3 \times 10^{5} \frac{\mathrm{~J}}{\mathrm{~kg}} \approx 3 \times 10^{8} \frac{\mathrm{~J}}{\mathrm{~m}^{3}}$

How fast does ice melt?
albedo $\alpha=0.7$ but less that far north, say $1000 \frac{\mathrm{~J}}{m^{2} S}$
rate of melting

$$
\frac{(1-0.7) 1000 \frac{\mathrm{~J}}{\mathrm{~m}^{2} s}}{3 \times 10^{8} \frac{\mathrm{~J}}{\mathrm{~m}^{3}}}=1 \times 10^{-6} \frac{\mathrm{~m}}{\mathrm{~s}}=9 \frac{\mathrm{~cm}}{\text { day }}
$$

# What do we want to know? 

melting of glacier<br>flow velocity of glacier<br>shear stress on base<br>effect of temperature on flow<br>shape of glacier

## What do we want to know?

## melting of glacier flow velocity of glacier shear stress on base effect of temperature on flow shape of glacier

## Part 1

## flow of viscous fluid between plates

## coordinate system



dynamical quantities

shear stress - strain rate law

$$
\sigma=\mu \frac{d v}{d z}
$$

Newton's Law

$$
\mu \frac{d^{2} v}{d z^{2}}+f=\rho \frac{d v}{d t}
$$

welded boundary
$v$ same as object
free boundary

$$
\sigma=0
$$

Newton's Law

$$
\begin{aligned}
& \mu \frac{d^{2} v}{d z^{2}}+\chi=\rho \frac{d v}{d t} \\
& \frac{d^{2} v}{d z^{2}}=0 \\
& \\
& \quad v(z=0)=v_{0} \\
& \\
& v(z=H)=0
\end{aligned}
$$

Newton's Law



## Lessons



## Part 2

flow of viscous fluid in a wide stream

## stream



## stream

no shear stress


$$
f=\sigma g \sin \theta
$$

## stream

Newton's Law


$$
f=\rho g \sin \theta
$$

$$
\begin{gathered}
\mu \frac{d^{2} v}{d z^{2}}+f=\rho \frac{d v}{d t} \\
\frac{d^{2} v}{d z^{2}}=-\frac{f}{\mu}=-B \\
\frac{d v}{d z}(z=0)=0 \\
v(z=H)=0
\end{gathered}
$$

## stream

Newton's Law


## stream

Newton's Law


## stream

Newton's Law


## stream

## Newton's Law

| no shear stress $\quad 0 \quad v$ | $A \frac{d^{2} v}{d z^{2}}=-B$ |
| :---: | :---: |
|  | $\text { В } \frac{d v}{d z}(z=0)=0$ |
| $\underset{\text { stationarystreambed }}{H} \underset{Z}{\square}$ | C $v(z=H)=0$ |
| $\begin{aligned} & v(z)=c_{0}+c_{1}(H-z)+c_{2}(H-z)^{2} \\ & d v / d z=-c_{1}-2 c_{2}(H-z) \\ & d^{2} v / d z^{2}=2 c_{2} \end{aligned}$ | $\xrightarrow{\text { males }} \Rightarrow c_{0}=0$ |

## stream

## Newton's Law



## stream

## Newton's Law



## stream

Newton's Law


$$
v(z)=B H(H-z)-1 / 2 B(H-z)^{2}
$$

## stream

## maximum velocity

$$
\begin{aligned}
& \text { no shear stress } \\
& \text { stationary streambed } \\
& v(z)=B H(H-z)-1 / 2 B(H-z)^{2} \\
& \sigma(z)=-\mu B H+\mu B(H-z)
\end{aligned}
$$

$$
v_{0}=v(0)=1 / 2 B H^{2}
$$

## maximum shear stress

$$
\begin{aligned}
\sigma_{H}=\sigma(H) & =-\mu B H \\
& =-2 \frac{\mu v_{0}}{H}
\end{aligned}
$$

$$
\sigma_{H}=\sigma(H)=-2 \frac{\mu v_{0}}{H}
$$

$$
\begin{aligned}
& H=1000 \mathrm{~m} \\
& v_{0}=1 \times 10^{-5} \mathrm{~m} / \mathrm{s} \quad(\sim \text { one meter per day }) \\
& \mu=10^{12} \mathrm{~Pa}-\mathrm{s} \\
& \sigma_{H}=-2 \frac{\mu v_{0}}{H}=-2 \times 10^{4} \mathrm{~Pa}=-20 \mathrm{kPa}
\end{aligned}
$$

similar to the strength of sands and gravels

# Part 3 <br> Glaciers <br> hotter at the bottom 

effect of variable (temperature dependent) viscosity


Newton's Law
Viscous Flow Law

$$
\frac{d \sigma}{d z}+f=\rho \frac{d v}{d t}
$$

$$
\sigma=\mu(z) \frac{d v}{d z}
$$

so using
chain rule

$$
\mu \frac{d^{2} v}{d z^{2}}+\frac{d \mu}{d z} \frac{d v}{d z}+f=\rho \frac{d v}{d t}
$$



Newton's Law

$$
\mu \frac{d^{2} v}{d z^{2}}+\frac{d \mu}{d z} \frac{d v}{d z}+\boldsymbol{X}=\rho \frac{d v}{d t}
$$



Newton's Law
$\mu \frac{d^{2} v}{d z^{2}}+\frac{d \mu}{d z} \frac{d v}{d z}=0 \quad \mu_{0} \exp (-c z) \frac{d^{2} v}{d z^{2}}-c \mu_{0} \exp (-c z) \frac{d v}{d z}=0$


Newton's Law

$$
\mu \frac{d^{2} v}{d z^{2}}+\frac{d \mu}{d z} \frac{d v}{d z}=0
$$

$$
\frac{d^{2} v}{d z^{2}}-c \frac{d v}{d z}=0
$$

$$
\begin{aligned}
& \frac{d^{2} v}{d z^{2}}-c \frac{d v}{d z}=0 \\
& \quad \text { let } Z=\frac{d v}{d z} \\
& \frac{d Z}{d z}-c Z=0 \quad \text { so } \quad \frac{d Z}{d z}=c Z \quad \text { so } \quad Z=Z_{0} \exp (c z) \\
& Z=\frac{d v}{d z} \quad \text { so } \quad v=\frac{Z_{0}}{c} \exp (c z)-C
\end{aligned}
$$

$$
v=\frac{Z_{0}}{c} \exp (c z)-C
$$

$$
\begin{aligned}
& \text { A } v(z=0)=v_{0} \\
& \text { В } v(z=H)=0
\end{aligned}
$$

B $\quad v=\frac{Z_{0}}{c} \exp (c H)-C=0 \quad$ so $\quad C=\frac{Z_{0}}{c} \exp (c H)$
A $\quad v=\frac{Z_{0}}{c}[\exp (0)-\exp (c H)]=v_{0}$

$$
\text { so } \quad Z_{0}=\frac{c v_{0}}{[1-\exp (c H)]}
$$

$$
v=v_{0} \frac{[\exp (c z)-\exp (c H)]}{[1-\exp (c H)]}
$$



## What about the shear stress

Newton's Law

$$
\begin{aligned}
& \frac{d \sigma}{d z}+\hat{m}=\rho \stackrel{d v}{d t} \\
& \\
& \text { so } \frac{d \sigma}{d z}=0
\end{aligned}
$$

and $\sigma=$ constant


small c
$\sigma=\mu \frac{d v}{d x}=\frac{\mu_{0} c v_{0}}{\left[1-\left(1+c H+1 / 2 c^{2} H^{2}\right)\right]}=\frac{\mu_{0} v_{0}}{H\left(1+1 / 2 c^{2} H\right)}$
stress less than constant viscosity case


Top part has more uniform velocity
Lessons
Stress went down (compared to uniform viscosity)

