

# Response of Atmospheric Convection to Vertical Wind Shear: Cloud Resolving Simulations with Parameterized Large-Scale Circulation.

## Part II: Effect of Interactive Radiation

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## ABSTRACT

The authors investigate the effects of cloud-radiation interaction and vertical wind shear on convective ensembles interacting with large scale dynamics in cloud resolving model simulations, with the large-scale circulation parameterized using the weak temperature gradient approximation. Numerical experiments with interactive radiation are conducted with imposed surface heat fluxes constant in space and time, an idealized lower boundary condition which prevents wind-evaporation feedback. Each simulation with interactive radiation is compared to a simulation in which the radiative heating profile is held constant in the horizontal and time, and is equal to the horizontal mean profile from the interactive radiation simulation with the same vertical shear and surface fluxes. Interactive radiation is found to reduce mean precipitation in all cases. The magnitude of the reduction is nearly independent of the vertical wind shear, but increases with surface fluxes. Deep shear also reduces precipitation, though by approximately the same amount with or without interactive radiation. The reductions in precipitation due to either interactive radiation or deep shear are associated with large-scale descent in the lower troposphere, which more strongly exports moist static energy and is quantified by a larger normalized gross moist stability.

## 20 **1. Introduction**

21 Radiative feedbacks have been shown to contribute to the dynamics of tropical  
22 disturbances on a range of time scales. The phenomena influenced by these feedbacks  
23 include tropical cyclogenesis (e.g. Bu et al 2014; Hakim 2013) and intraseasonal  
24 variability (e.g., Raymond 2001; Bony and Emanuel 2005; Kim et al. 2011; Sobel and  
25 Maloney 2013), and have been found in observations (e.g. Lin and Mapes 2004; Sobel et  
26 al. 2014) as well as idealized numerical modeling studies (e.g., Raymond 2001;  
27 Grabowski and Moncrieff 2002; Bony and Emanuel 2005; Wing and Emanuel 2014).

28 By absorbing longwave radiation and re-emitting it to space at lower temperatures,  
29 clouds and water vapor warm the troposphere. Locally, regions of deep convection  
30 experience a strong reduction in the net radiative cooling compared to adjacent clear-sky  
31 regions. At the top of the atmosphere, this longwave warming is often largely (though  
32 imperfectly; Lin and Mapes 2004), offset by shortwave anomalies due to the enhanced  
33 albedo of deep clouds (e.g., Ramanathan et al. 1989; Hartmann et al. 2001), so that the  
34 effect of high tropical clouds on the long-term climate is modest. However, the longwave  
35 effect is felt immediately in the atmosphere, while the compensating shortwave effect is  
36 felt largely at the surface. If the surface is ocean, the heat capacity of the ocean mixed  
37 layer can delay the surface response enough that on short time scales, the positive  
38 feedback in the atmosphere dominates. By acting as a source of column-integrated moist  
39 static energy where that quantity already has a positive anomaly, this longwave radiative  
40 feedback is positive, and can organize tropical convection on a range of scales, including  
41 the intraseasonal (Raymond 2001; Sobel and Gildor 2003; Grabowski and Moncrieff

42 2004; Bretherton et al. 2005; Wing and Emanuel 2014; Emanuel et al. 2014; Tobin et al.  
43 2013; Chikira 2014).

44 We expect the effectiveness of cloud-radiative feedbacks to be related to the  
45 horizontal extent, vertical structure and microphysical properties of the cloud field. These  
46 cloud properties, in turn, are influenced by a number of environmental factors. Vertical  
47 wind shear is an important such factor, in that we expect it to play a significant role in the  
48 organization of convective cloud systems on the mesoscale. In this study, we describe  
49 studies designed to isolate the effects of interactive radiation and vertical shear, together  
50 and separately, on tropical oceanic convection.

51 In the first part of this study (Anber et al. 2014 hereafter AWS14), we performed  
52 CRM simulations in which we prescribed both surface fluxes and column radiative  
53 heating, and introduced a background vertical wind shear of different strengths as a  
54 control parameter. We found that under otherwise identical large-scale conditions, the  
55 precipitation can change significantly as a result of the wind shear influence, as the shear  
56 organizes the convective systems into squall line-like structures. These structures  
57 influence the convective heating and large-scale vertical motion profiles, with the  
58 relationship between heating and large-scale vertical motion parameterized in AWS14 by  
59 the weak temperature gradient (WTG) approximation. The different vertical motion  
60 profiles lead to different mean precipitation rates for the same export of moist static  
61 energy (with that export fixed in the experiments of AWS14 by the imposed surface  
62 fluxes and radiative cooling), a relationship quantified by the normalized gross moist  
63 stability.

64           In this sequel paper we perform similar model integrations, again studying the  
65 influence of vertical wind shear on tropical convection, but now with interactive radiative  
66 heating. The results are then compared to those obtained from otherwise identical  
67 experiments in which the radiative heating is prescribed to have the same profile in the  
68 time mean as the interactive radiation simulations.

69           Our simulation design is such that we are not studying the most direct aspects of  
70 the radiative feedback on convective organization, in which the radiative heating itself  
71 changes on the large scale (represented in our simulations by the horizontal and time  
72 mean over the domain). We compare interactive- and fixed-radiation simulations in  
73 which the radiative heating itself is the same in the horizontal and time mean. The direct  
74 radiative feedback relevant to tropical disturbances can be estimated straightforwardly,  
75 from observations or global model simulations, as the column-integrated radiative  
76 heating anomaly per unit anomaly in column water vapor, precipitation, or other metric  
77 of convective activity on the space and time scale of interest (e.g., Bretherton and Sobel  
78 2002; Su and Neelin 2002; Lin and Mapes 2004). Rather, we investigate the more  
79 indirect effects which result not from the change in the areally averaged radiative heating,  
80 but from the effects that *interactive* radiation may have on cloud structure and dynamics,  
81 holding the large-scale mean radiative heating profile fixed. These interactivity effects  
82 then change the cloud system structure and horizontal mean latent heating, and influence  
83 the export of moist static energy via changes in the large-scale vertical motion profile  
84 associated with those latent heating changes. We find that these effects can be quite  
85 significant. We are not aware of any previous documentation of them.

86 To facilitate our focus on the vertical shear influence and cloud-radiative  
87 interaction, we prescribe surface turbulent heat fluxes. This eliminates surface flux  
88 feedbacks which would occur if the lower boundary were taken to be at fixed sea surface  
89 temperature (SST). With varying vertical shear, changes in momentum transport to the  
90 surface can change surface winds from one simulation to the next, inducing strong  
91 surface turbulent flux feedbacks under fixed SST. These feedbacks are not of interest  
92 here, so we eliminate them by prescribing surface turbulent heat fluxes directly. We have  
93 also conducted simulations with a slab ocean model; these increase the size of the  
94 parameter space and the range of possible model behavior enough that we will report on  
95 them in a separate study.

96 After describing the model and experiment design in section 2, we present the  
97 results in section 3. We summarize in section 4.

98

## 99 **2. Model and Experiment Design**

100 The model used here is the Weather Research and Forecast (WRF) model version  
101 3.3, in three spatial dimensions with the same physical parameterization schemes used in  
102 AWS14. For experimental design and implementation of the weak temperature gradient  
103 (WTG) approximation (Sobel and Bretherton 2000; Raymond and Zeng 2005) the reader  
104 is referred to AWS14. All experiments are integrated more than 40 days to achieve  
105 statistical equilibrium, and analysis of the output is conducted on the last 20 days, during  
106 which the simulations have reached statistical equilibrium.

107 To obtain the reference temperature profile required for WTG, we follow a  
108 standard procedure (e.g., Sobel and Bretherton 2000; Raymond and Zeng 2005; Wang

109 and Sobel 2011; Daleu et al. 2015). We first perform a radiative-convective equilibrium  
110 (RCE) experiment at a fixed SST of 28 °C with interactive radiation, reaching equilibrium  
111 at about 60 days. Results from this experiment are averaged over the last 10 days after  
112 equilibrium to obtain statistically equilibrated temperature and moisture profiles. Figure 1  
113 shows the resulting vertical profiles of (a) potential temperature and (b) moisture. These  
114 profiles are then used to initialize runs with parameterized large scale circulation, and the  
115 temperature profile is used as the target profile against which perturbations are computed  
116 in the WTG runs.

117 For our experiments with interactive radiation, the NCAR Community  
118 Atmosphere Model (CAM 3.0) shortwave and longwave radiation schemes (Collins et al.  
119 2004) are used. Diurnal and seasonal cycles are removed and the top of the atmosphere  
120 incident solar flux is set to  $400 \text{ W m}^{-2}$ .

121 For WTG runs without vertical wind shear, we prescribe total turbulent surface  
122 fluxes of heat and moisture at three different values: low ( $140 \text{ W m}^{-2}$ ), moderate ( $180 \text{ W}$   
123  $\text{m}^{-2}$ ), and high ( $220 \text{ W m}^{-2}$ ); these variations are achieved holding the sensible heat flux  
124 fixed at  $20 \text{ W m}^{-2}$  while the latent heat flux takes on values of 120, 160, and  $200 \text{ W m}^{-2}$ ,  
125 respectively.

126 We constrain the vertical wind shear by relaxing the horizontal mean zonal wind  
127 towards a prescribed profile with a time scale of one hour, as in AWS14. We vary the  
128 shear depth over three layers; shallow (1500 m), midlevel (4500 m), and deep (12000 m),  
129 while holding the target wind speed at the top of the shear layer constant at 20 m/s, and  
130 keeping it zero at the surface (see Fig.1 of AWS14). Surface fluxes for this set of  
131 experiments are held fixed at 140, 180, and  $220 \text{ W m}^{-2}$  but we show results from the case

132 of  $180 \text{ W m}^{-2}$  only since results from other cases are similar and same conclusion can be  
133 made. Clouds and water vapor then evolve freely and strongly influence the radiative  
134 heating profile.

135 The time-and domain-averaged radiative heating profiles for this set of  
136 experiments is shown in Fig. 2*a* for the unsheared environment with different surface  
137 fluxes, and Fig. 2*b* for different shear depths with intermediate surface fluxes. For either  
138 low surface fluxes or shallow shear there is a cooling throughout the whole atmospheric  
139 column that reaches 1 K/day in the upper troposphere. For higher surface fluxes or deeper  
140 shear, there is a significant radiative heating in the mid-to upper troposphere, exceeding 5  
141 K/day for high surface fluxes. Above 10 Km there is a sharp drop in the radiative heating  
142 profile due to the strong emission of longwave radiation from cloud top, with radiative  
143 cooling exceeding 10 K/day for strong surface fluxes or deep shear.

144 To compare the effect of interactive radiation on the cumulus ensemble  
145 interacting with the parameterized large scale dynamics, we perform a set of experiments  
146 identical to those described above except that we prescribe the radiative cooling profile in  
147 Figure 2 and keep it constant in both the horizontal and time. This the time-and domain-  
148 averaged radiative cooling is the same with either fixed or interactive radiation. As  
149 surface fluxes are already fixed, this means that both the equilibrated time mean column-  
150 integrated moist static surface fluxes and net radiative cooling are identical in pairs of  
151 interactive- and fixed-radiation experiments. Any differences in the time mean column  
152 moist static energy budgets must be due to the export term, which is controlled by the  
153 large-scale vertical velocity profile.

154

### 155 3. Results

156 Fig.3 shows time and domain averaged precipitation,  $P$ , as a function of (a)  
 157 surface fluxes in unsheared environment and (b) wind shear at moderate surface fluxes of  
 158  $180 \text{ W m}^{-2}$  (shear cases at low and high surface fluxes are similar and not shown), for  
 159 experiments with interactive radiation (blue) and prescribed radiation (red). In all pairs of  
 160 experiments, interactive radiation produces significantly less precipitation than does  
 161 prescribed radiation. The difference monotonically increases with surface fluxes (Fig.  
 162 3a). The difference is about 1mm/d for small surface fluxes, and reaches 3 mm/d and 5  
 163 mm/d for the moderate and high surface fluxes, respectively, while remaining nearly  
 164 constant with shear (Fig. 3b). Mid-level shear produces the greatest precipitation in both  
 165 interactive as well as prescribed radiation experiments. A non-monotonic relation  
 166 between shear and precipitation was also found in the first part of this study, with  
 167 specified radiative cooling, and in other studies with different methodologies (e.g.  
 168 AWS14, Weisman et al. 2004).

169 Precipitation is also diagnosed from the combined vertically integrated budgets of  
 170 moist and dry static energy as:

$$171 \quad P = \frac{1}{M} (L + H + \langle Q_R \rangle) - \langle Q_R \rangle - H$$

172 where  $L, H, \langle Q_R \rangle$ , and  $M$  are the latent heat flux, sensible heat flux, vertically  
 173 integrated radiative heating, and normalized gross moist stability, respectively. The latter

$$174 \quad \text{is defined as } M = \left\langle \overline{W} \frac{\partial h}{\partial z} \right\rangle / \left\langle \overline{W} \frac{\partial s}{\partial z} \right\rangle$$

175 where  $W, s$ , and  $h$  are the large-scale vertical velocity, dry, and moist static energy (MSE)  
 176 respectively, brackets denote mass weighted vertical integrals, and overbars are domain

177 and time mean. The above equation gives very good agreement with the model output as  
178 indicated in Fig.3 with square symbols. The above equation helps us to interpret changes  
179 in  $P$ .

180 To understand the behavior of precipitation in the presence of cloud-radiation  
181 interaction, we look at the large-scale vertical motion  $W_{WTG}$ . Figure 4 shows the vertical  
182 profile of  $W_{WTG}$  with surface fluxes of 140, 180, and 220  $\text{W m}^{-2}$  ((a), (b), and (c),  
183 respectively), in the unsheared environment, and shear depths of 1500 m, 4500 m, and  
184 12000 m ((d), (e), and (f), respectively), with intermediate surface fluxes, for both  
185 interactive (solid curves) and prescribed (dashed) radiative cooling. In all cases, when  
186 radiation is interactive the large scale vertical motion exhibits a qualitatively similar  
187 vertical structure to that when radiation is prescribed, but with a sharper peak at about 10  
188 km and stronger subsidence in the lower levels (except in the shallow shear case, Fig.  
189 4d), especially for the high surface flux and deep shear cases.

190 A closer look at the peak magnitude of the ascent and descent is shown in Fig. 5,  
191 which displays the maximum and minimum values of the  $W_{WTG}$  profile as a function of  
192 surface fluxes and shear depth. The peak amplitude of the vertical motion is a  
193 monotonically increasing function of the surface fluxes and the shear depth for both  
194 interactive and prescribed radiative heating. However, there are significant differences  
195 between the two mean profiles. Except for the shallow shear depth where the heating  
196 anomalies are negative, such heating anomalies increase with either increasing surface  
197 fluxes or shear depth, and reach maximum amplitude for the highest surface fluxes and  
198 largest shear depth.

199

200           The most striking feature in the large scale vertical velocity structure is the  
201 significant subsidence in the lower troposphere, resulting from evaporative cooling under  
202 mesoscale cloud systems. For low and moderate surface fluxes, interactive radiation  
203 introduces a small increase in this subsidence, and for shallow and midlevel shear the  
204 difference in subsidence introduced by interactive radiation is small. However, for the  
205 high surface fluxes and deep shear cases, interactive radiation strongly strengthens this  
206 lower-tropospheric subsidence. This is particularly important in the moist static energy  
207 budget.

208

209           This indirect role of interactive radiation in the thermodynamics can be usefully  
210 interpreted in terms of the normalized gross moist stability (e.g. Neelin and Held 1987;  
211 Sobel 2007) shown in Figure 6, and defined above as the export of moist static energy  
212 normalized by the export of dry static energy. Increasing surface fluxes while holding the  
213 radiative heating prescribed does not have a strong impact on  $M$ . For interactive  
214 radiation, however, as the surface fluxes increase,  $M$  increases strongly. Introduction of  
215 deep shear also increases  $M$ , but does so similarly whether radiation is fixed or  
216 interactive. These increases are results of stronger ascent in the upper troposphere and  
217 descent in the lower troposphere, implying horizontal divergence both at upper levels and  
218 near the surface where moist static energy is large, and convergence in mid-troposphere  
219 where it is smaller.  $M$  exceeds unity for the greatest surface fluxes and shear depth,  
220 meaning that the export of moist static energy exceeds the export of dry static energy, as  
221 seen late in the active phases of the MJO (e.g. Wang et al. 2015; Sobel et al. 2014). Any  
222 increase in radiative heating or surface fluxes for  $M > 1$  is relatively inefficient in

223 increasing precipitation since  $P$  is inversely proportional to the product of  $1/M$  and these  
224 MSE sources.

225         The strong increase in  $M$  in the interactive vs. fixed radiation experiments was  
226 unexpected, and implies some cancellation between direct and indirect effects of  
227 radiation in the moist static energy budgets of deep convective tropical disturbances. In  
228 our simulations, the direct radiative feedback (defined operationally here as the horizontal  
229 mean radiative cooling change as a function of precipitation or other measure of  
230 convective activity) is present only implicitly, since interactive- and fixed-radiation  
231 experiments with the same surface fluxes share the same time-mean radiative heating  
232 profiles. We expect this direct feedback to be positive, reducing the radiative cooling in  
233 the presence of large high cloud cover and column humidity and thus providing an  
234 anomalous source of MSE to the column. Our results imply that the radiative feedback  
235 also has a significant indirect effect which opposes this, changing the large-scale vertical  
236 motion profile so as to drive stronger advective export of MSE, thus acting as an  
237 anomalous sink of that quantity.

238 This result may be relevant to a number of observed phenomena; one that occurs to us in  
239 particular is the MJO. As shown in several studies (e.g., Yasunaga and Mapes 2012; Kim  
240 et al. 2014; Sobel et al. 2014; Inoue and back 2015a, 2015b), the maximum in  
241 precipitation associated with the MJO active phase occurs nearly simultaneously with the  
242 maximum in column-integrated MSE over the Indian Ocean. Column-integrated vertical  
243 advection of MSE is negative (i.e., export; positive  $M$ ) during the active phase, and  
244 increases in magnitude as the phase advances, with minimum  $M$  before the peak MSE  
245 and precipitation and maximum after it (in either longitude or time). These variations in

246 vertical advection are associated with the MJO's eastward propagation, as anomalous  
247 MSE import ahead acts to increase MSE leading up to the peak in the active phase, and  
248 anomalous export afterwards acts to reduce it (e.g., Kiranmayi and Maloney 2011; Kim et  
249 al. 2014; Sobel et al 2014; Inoue and Back 2015a, b; Yokoi and Sobel 2015).

250         The progression in the vertical motion profile to greater top-heaviness as the  
251 MJO's active phase matures is associated with a gradual transition from more convective  
252 to more stratiform cloud types, similar to what happens in individual mesoscale systems,  
253 but on much larger scales in both space and time (Mapes 2006). While the dynamics of  
254 the convective-stratiform transition are understood on the mesoscale, it has been unclear  
255 what mechanisms cause such a similar transition to play out on the planetary scale of the  
256 MJO, where it is not associated with individual mesoscale convective systems but rather  
257 a much larger envelope containing many such systems. While by no means solving this  
258 puzzle completely, our results suggest that both vertical shear and cloud-radiative  
259 interactions – both of which are strong in the late active phase of the MJO - may play key  
260 roles on this larger scale.

261

262         Finally, the temperature anomalies associated with interactive radiation are  
263 relevant to understanding the changes in vertical motion described above. Figure 7 shows  
264 the mean temperature anomalies from the RCE temperature profile as a function of  
265 surface fluxes (a), and shear depths (b). The low surface flux and shallow shear cases  
266 show warm anomalies throughout the atmosphere. However, as the surface fluxes or  
267 shear depth increase, the warm anomalies in the upper troposphere, and cool anomalies in  
268 the lower troposphere both strengthen. These temperature anomaly profiles indicate a

269 more stable environment, as usually observed in mesoscale convective systems and  
270 associated with horizontally extensive stratiform cloud and a second baroclinic mode  
271 component to the heating profile (Houze 1989).

272

## 273 **5. Summary**

274 We have presented results from idealized 3D cloud resolving simulations with  
275 large-scale circulation parameterized under the weak temperature gradient approximation  
276 to investigate the effects of interactive radiation and vertical shear on tropical deep  
277 convection. Surface fluxes of heat and moisture were prescribed to remove wind-surface  
278 flux feedbacks, and results were presented for different imposed vertical wind shear  
279 profiles as well as different surface flux magnitudes. Results from simulations with fixed  
280 and interactive radiation were compared, with the two types of simulations having the  
281 same time-mean radiative heating profile so that any difference is only due to the  
282 interactivity itself. The direct radiative-convective feedback, associated with changes in  
283 the domain-averaged radiative heating profile (e.g., as a function of precipitation or  
284 column water vapor) was excluded.

285 Interactive radiation has a significant impact when high surface fluxes are high,  
286 by generating a stronger envelope of high anvil cloud that extends over a large horizontal  
287 extent. The heating profile associated with this cloud system, via WTG, strengthens  
288 large-scale ascent in the upper troposphere and causes descent in the lower troposphere.  
289 This results in an increase in the normalized gross moist stability, and a decrease in mean  
290 precipitation, as moist static energy is more efficiently exported by the more top-heavy  
291 ascent profile. Deep vertical shear has a similar effect; it has this effect with either fixed

292 or interactive radiation, but the reduction in precipitation due to strong moist static energy  
293 is greatest when both deep shear and interactive radiation are present. We suggest that  
294 this behavior is relevant to the late active phase of the Madden-Julian oscillation (MJO),  
295 in which deep, strong vertical shear is present and surface fluxes are large. Coupling  
296 between radiation and dynamics of the type simulated here may be relevant to the  
297 increase in normalized gross moist stability which is found in observations late in the  
298 active phase of the MJO (Kim et al. 2014; Sobel et al. 2014; Inoue and Back 2015a,b),  
299 and which plays a significant role in the MJO's eastward propagation.

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301

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490 vertical wind shear depths, for simulations of interactive (blue), and prescribed, non-  
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493 FIG.7. Domain mean profile of temperature anomalies, as departure from RCE profile, as

494 a function of (a) surface fluxes, and (b) vertical shear depths.

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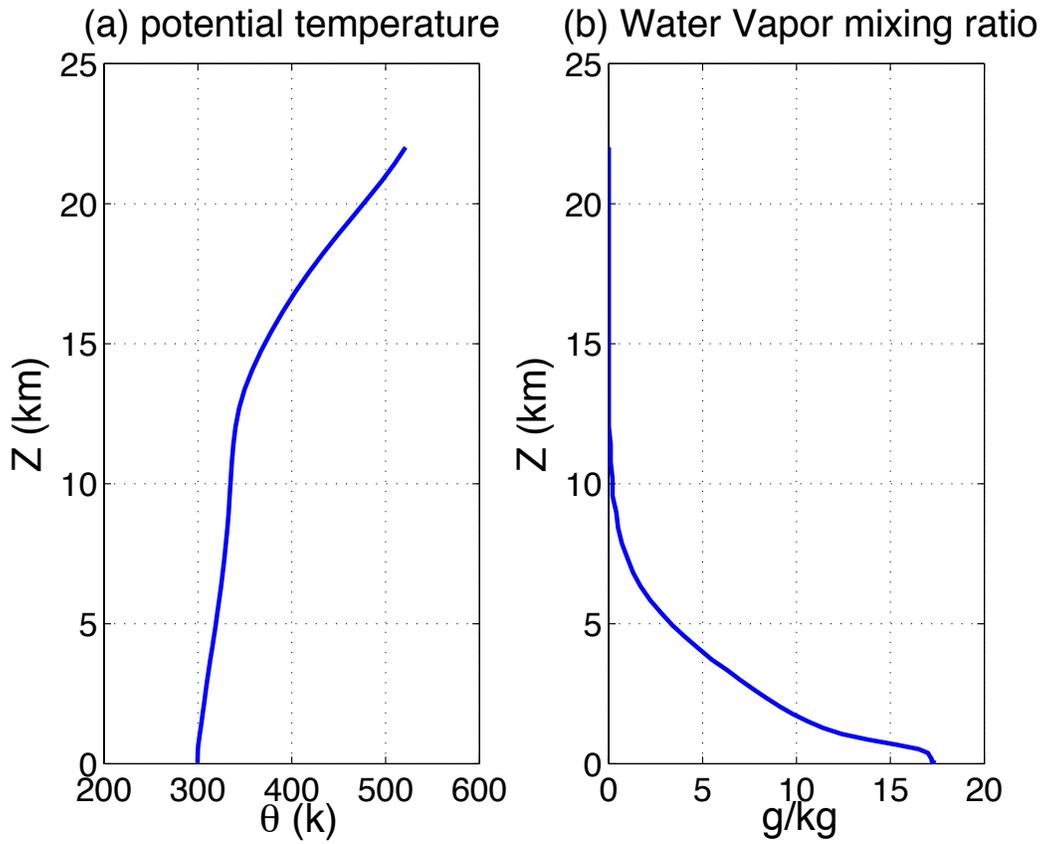
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515 **Figures**

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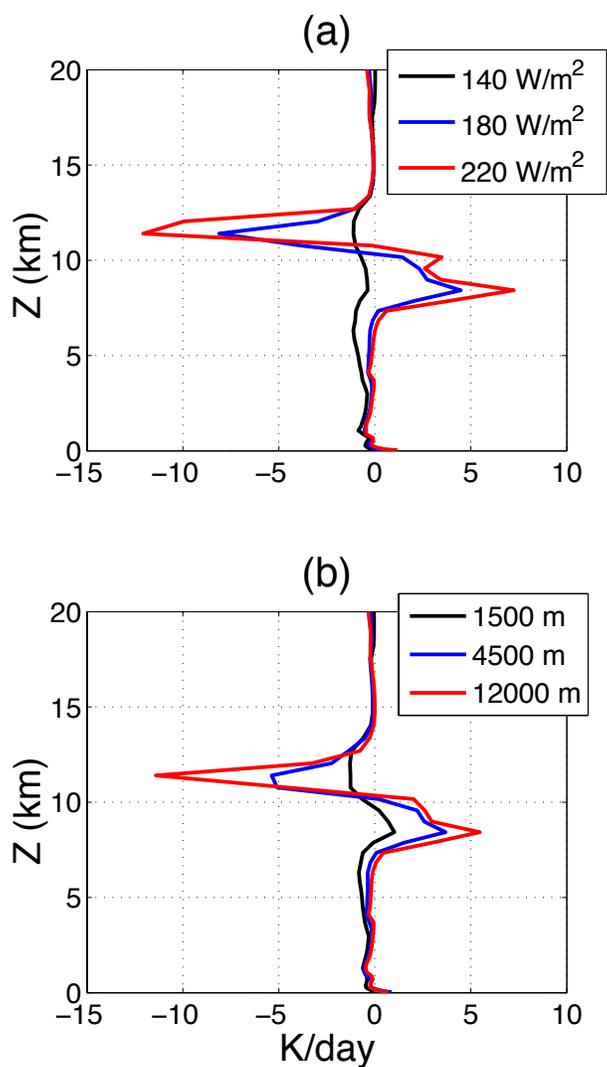
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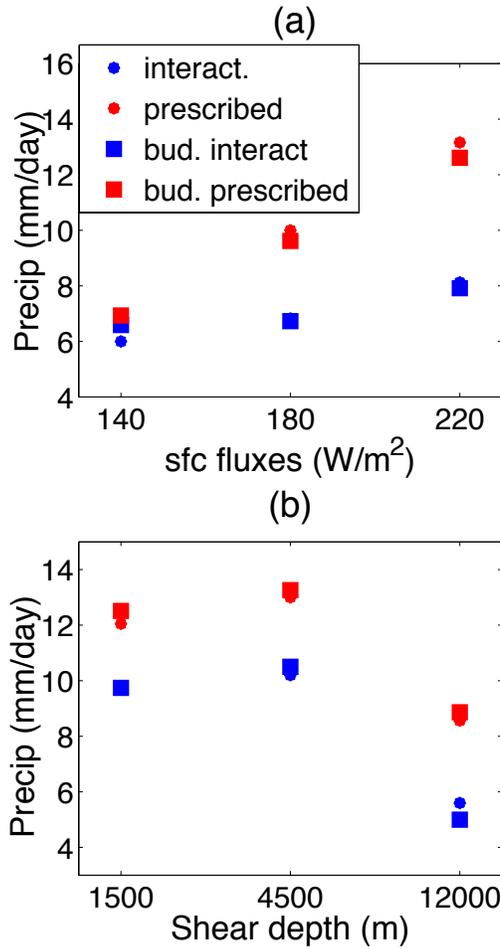
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525 FIG.2. Domain averaged radiative cooling profile as a function of (a) surface fluxes in  
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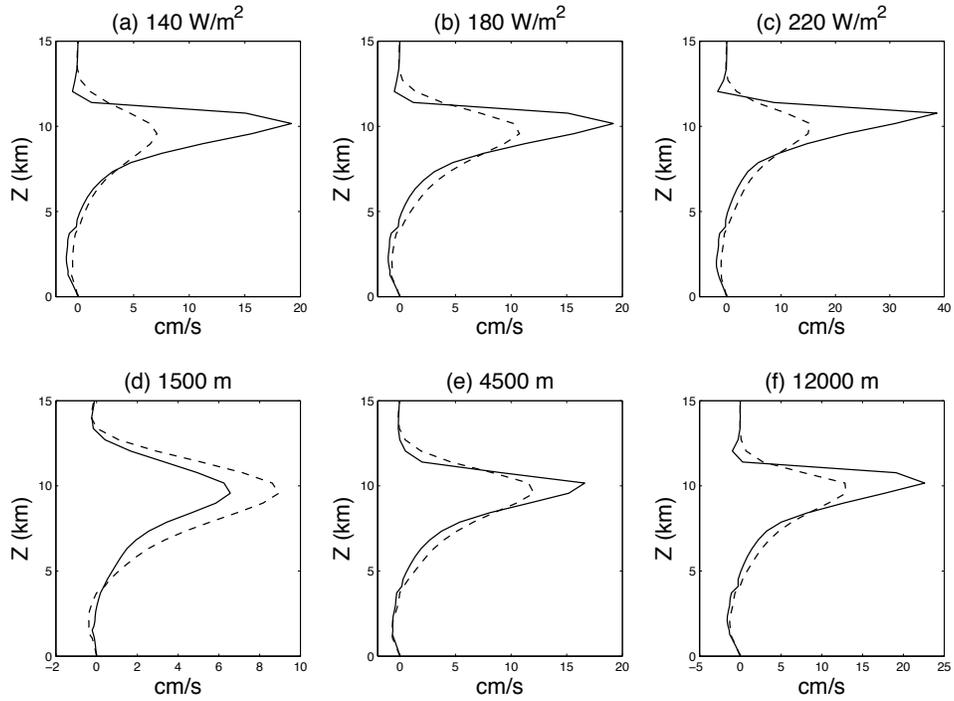
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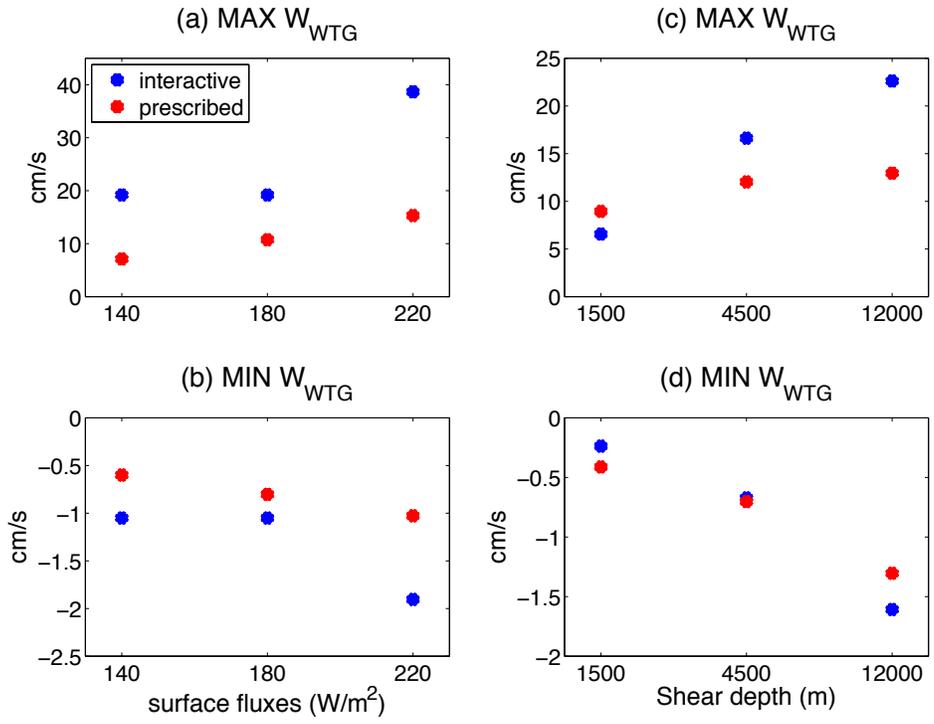
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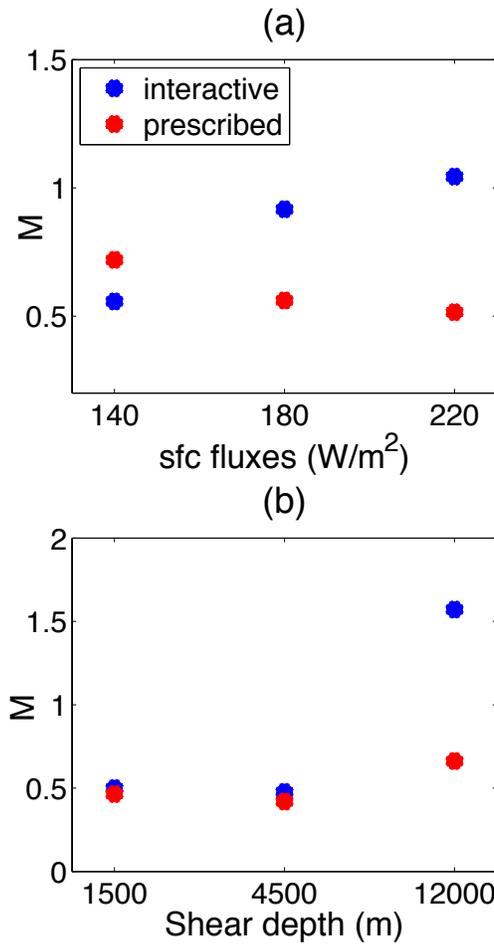
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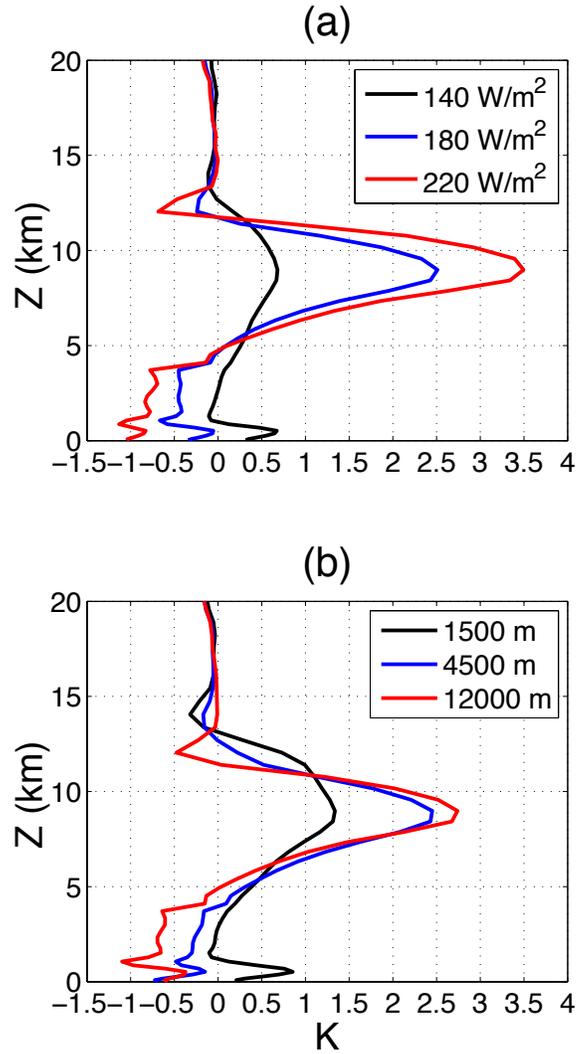
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