

Solid Earth Dynamics

Bill Menke, Instructor

Lecture 7

Today:

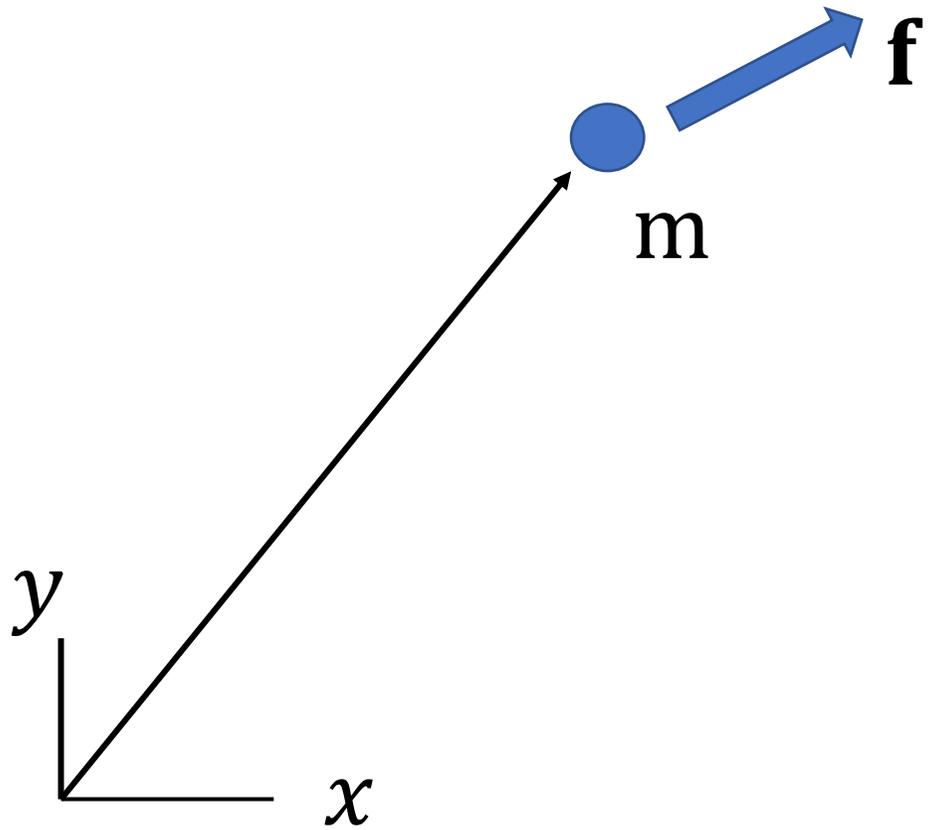
Gravity:

Newtonian orbits

field of a point mass and sphere

measuring gravity

ocean surface



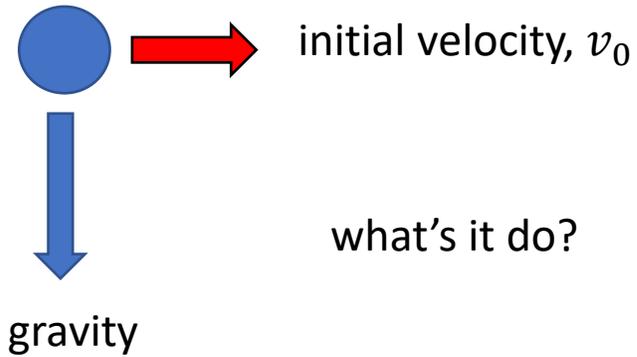
force = mass x acceleration

$$f_x = m \frac{d^2 x}{dt^2}$$

$$f_y = m \frac{d^2 y}{dt^2}$$

$$f_x = 0$$

$$f_y = -gm$$



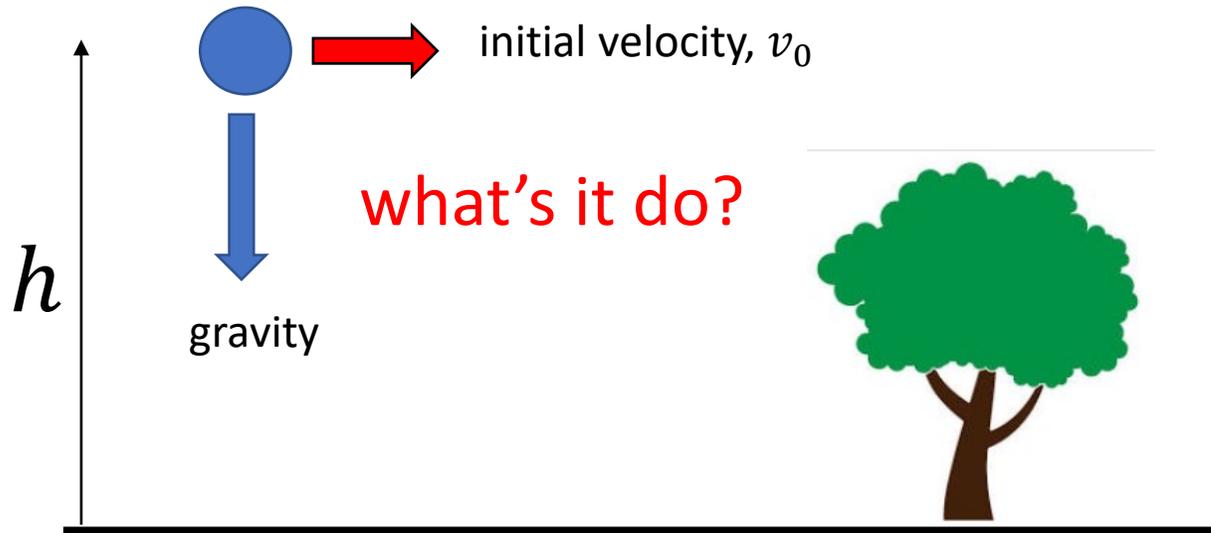
force = mass x acceleration

$$0 = m \frac{d^2 x}{dt^2}$$

$$-gm = m \frac{d^2 y}{dt^2}$$

$$f_x = 0$$

$$f_y = -mg$$



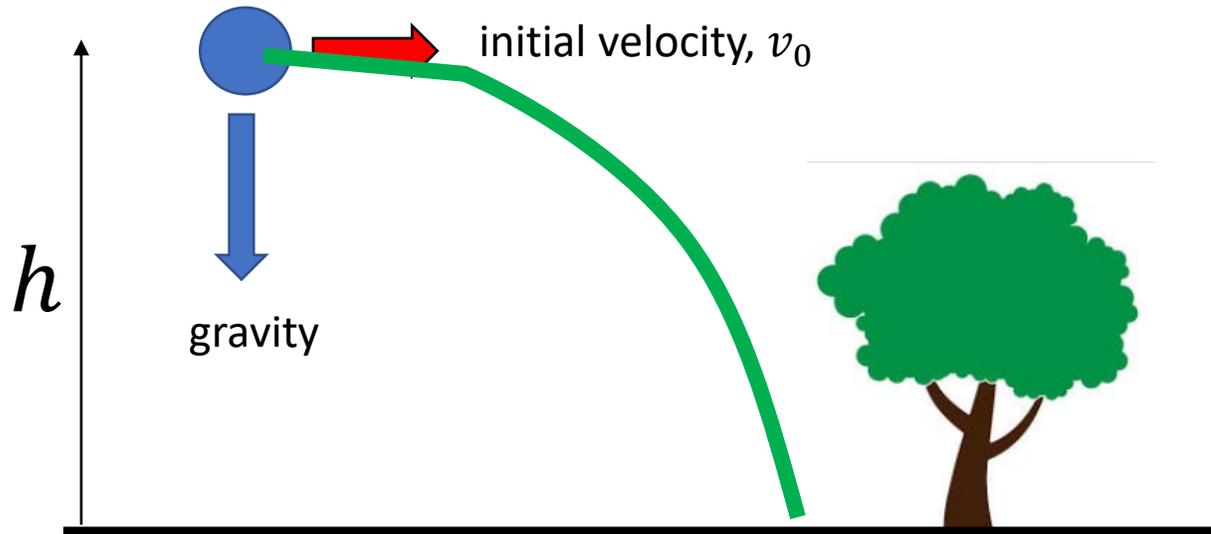
force = mass x acceleration

$$0 = \frac{d^2x}{dt^2}$$

$$-g = \frac{d^2y}{dt^2}$$

$$f_x = 0$$

$$f_y = -mg$$



force = mass x acceleration

$$0 = \frac{d^2x}{dt^2}$$

$$x = v_0 t$$

$$-g = \frac{d^2y}{dt^2}$$

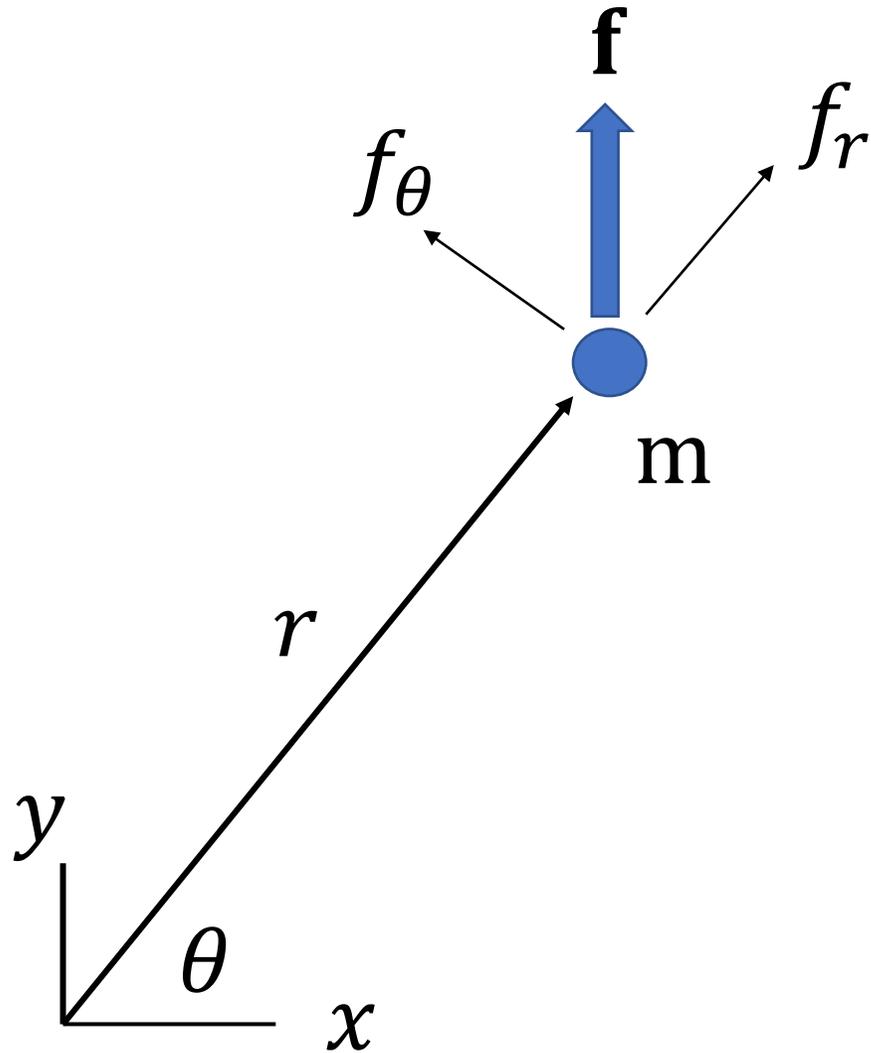
$$y = h - \frac{1}{2}gt^2$$

path doesn't depend on mass



really need to work in polar coordinates

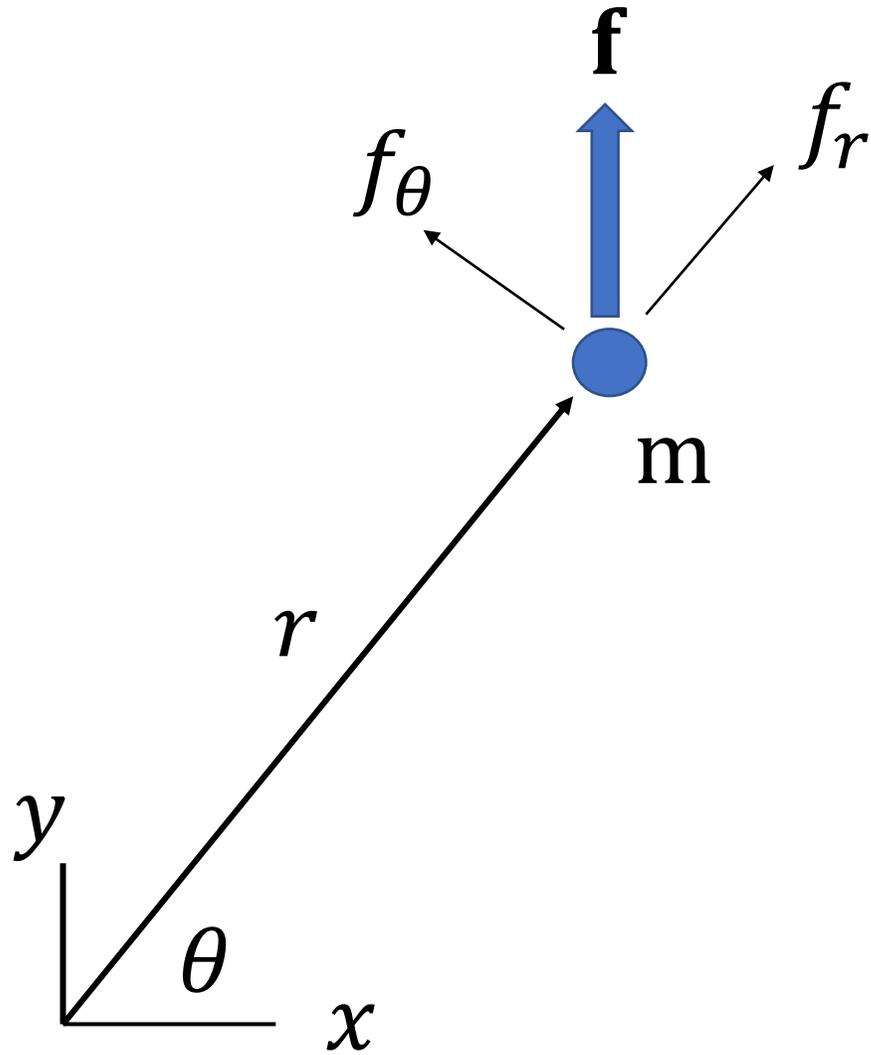




force = mass x acceleration

$$f_r = m ?$$

$$f_\theta = m ?$$



force = mass x acceleration

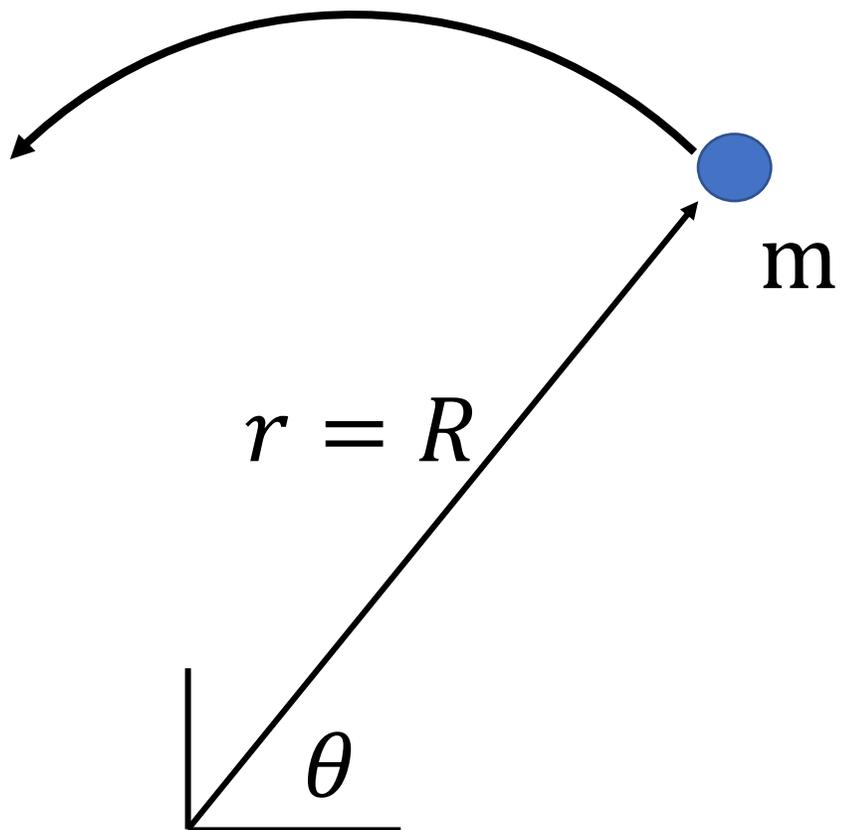
$$f_r = m \left(\frac{d^2 r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 \right)$$

$$f_\theta = m \left(r \frac{d^2 \theta}{dt^2} + 2 \frac{dr}{dt} \frac{d\theta}{dt} \right)$$

constant rate in circular orbit

$$r = R = \text{constant}$$

$$\theta = \frac{2\pi t}{T} \text{ with period, } T$$



$$f_r = m \left(\cancel{\frac{d^2 r}{dt^2}} - r \left(\frac{d\theta}{dt} \right)^2 \right)$$

$$f_\theta = m \left(r \cancel{\frac{d^2 \theta}{dt^2}} + 2 \frac{\cancel{dr}}{dt} \cancel{\frac{d\theta}{dt}} \right)$$

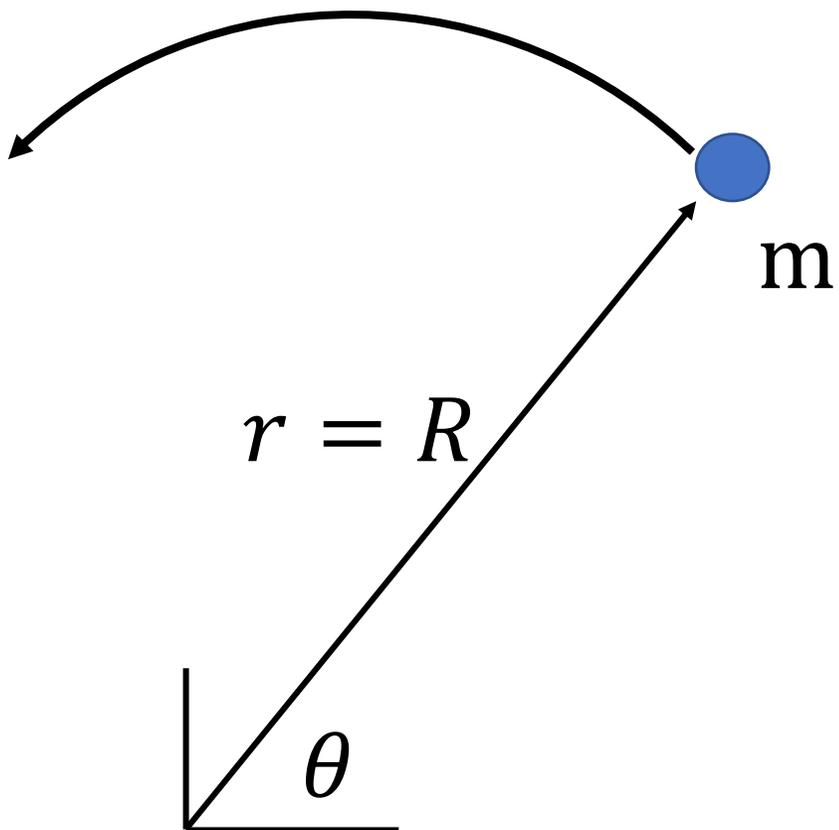
constant rate in circular orbit

$$r = R = \text{constant}$$

$$\theta = \frac{2\pi t}{T} \text{ with period, } T$$

$$f_r = -mR \left(\frac{2\pi}{T} \right)^2$$

$$f_\theta = 0$$



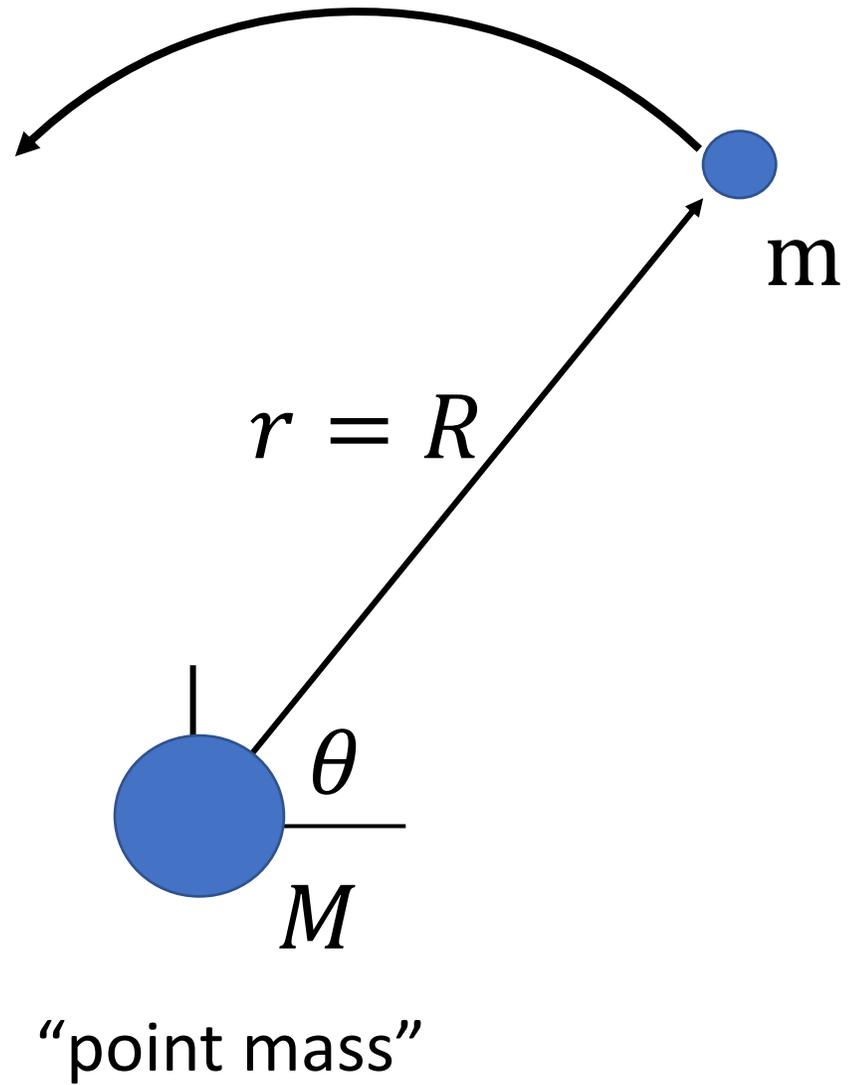
Kepler's "Law"

based on planetary observations

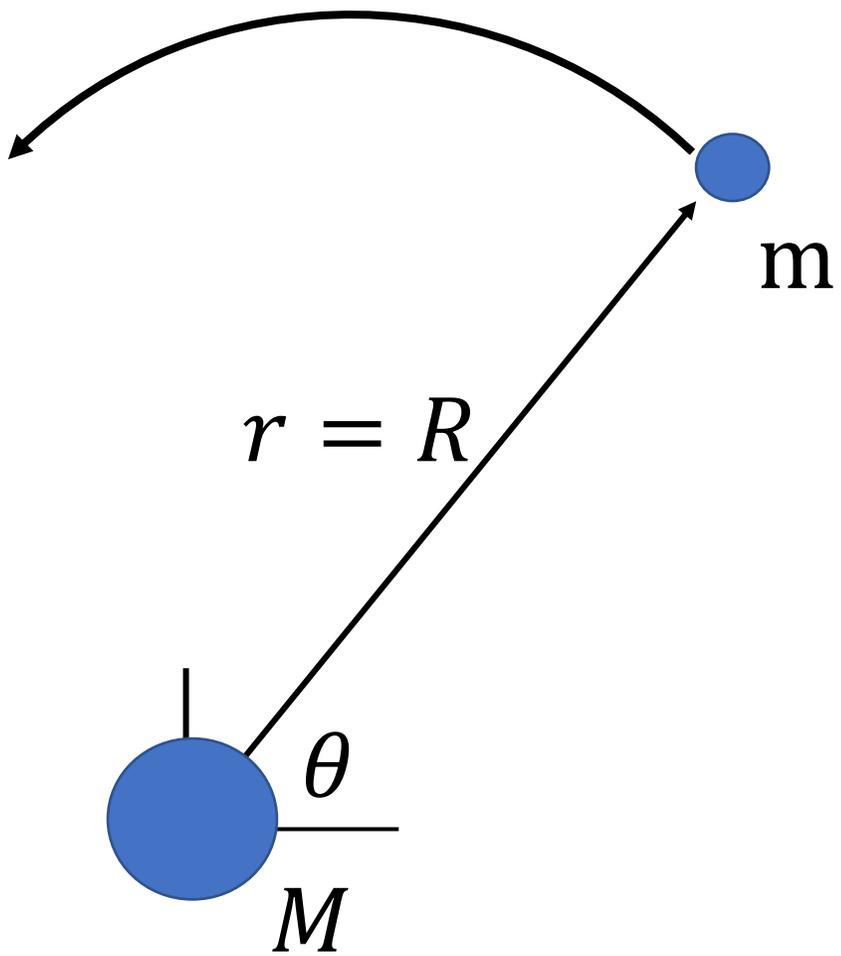
$$T^2 \propto R^3$$



insert Kepler's 3rd Law



$$\begin{aligned} f_r &= -mR \left(\frac{2\pi}{T} \right)^2 \\ &= -mR \frac{(2\pi)^2}{T^2} \\ &\propto -mR \frac{(2\pi)^2}{R^3} = -\frac{\gamma m M}{R^2} \end{aligned}$$



Newton's "Law" of gravity

$$f_{\theta} = 0$$

$$f_r = -mR \left(\frac{2\pi}{T} \right)^2$$

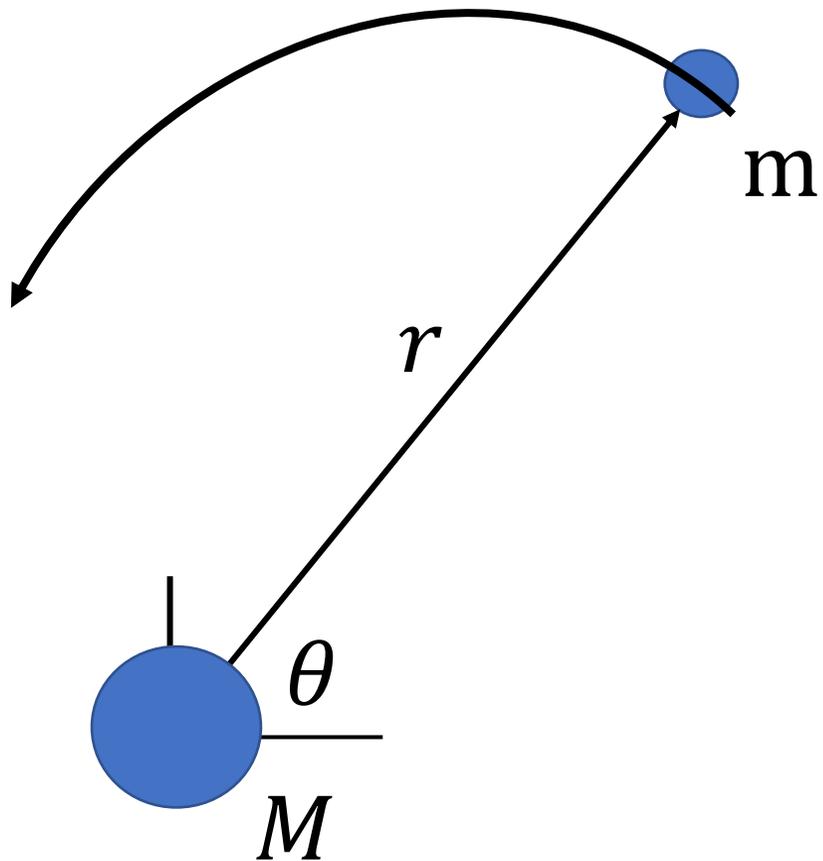
$$= -mR \frac{(2\pi)^2}{T^2}$$

$$\propto -mR \frac{(2\pi)^2}{R^3} = -\frac{\gamma m M}{R^2}$$



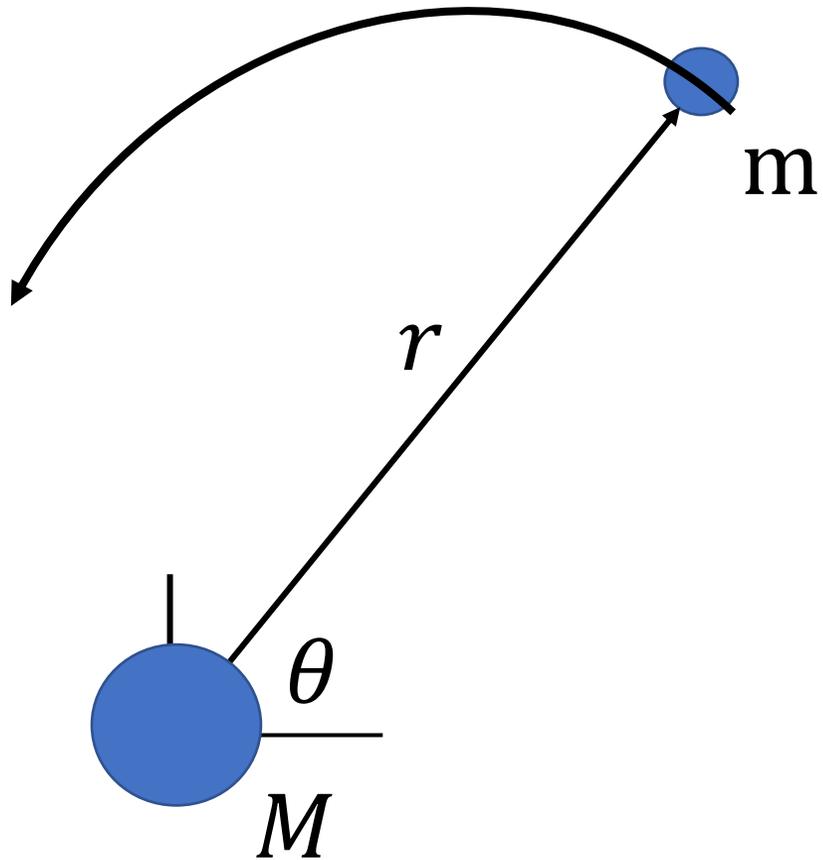
non-circular orbit

$$C = \gamma M$$



$$-\frac{Cm}{r^2} = m \left(\frac{d^2 r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 \right)$$

$$0 = m \left(r \frac{d^2 \theta}{dt^2} + 2 \frac{dr}{dt} \frac{d\theta}{dt} \right)$$

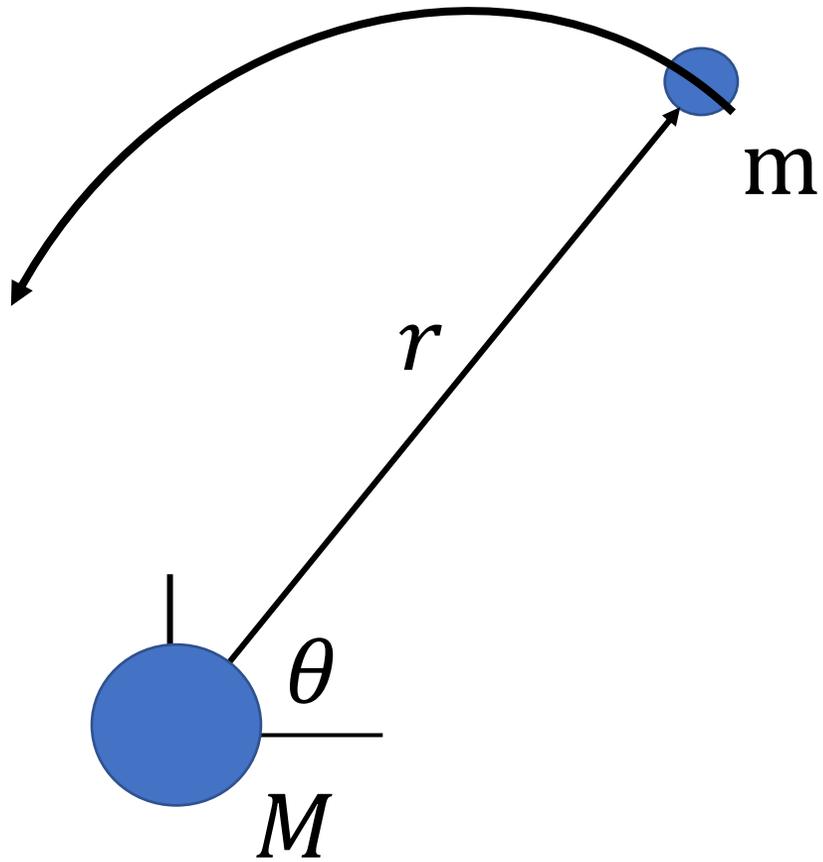


chain rule

$$\frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = r^2 \frac{d^2\theta}{dt^2} + 2r \frac{dr}{dt} \frac{d\theta}{dt}$$

multiply by r

$$r^2 \frac{d^2\theta}{dt^2} + 2r \frac{dr}{dt} \frac{d\theta}{dt} = 0$$

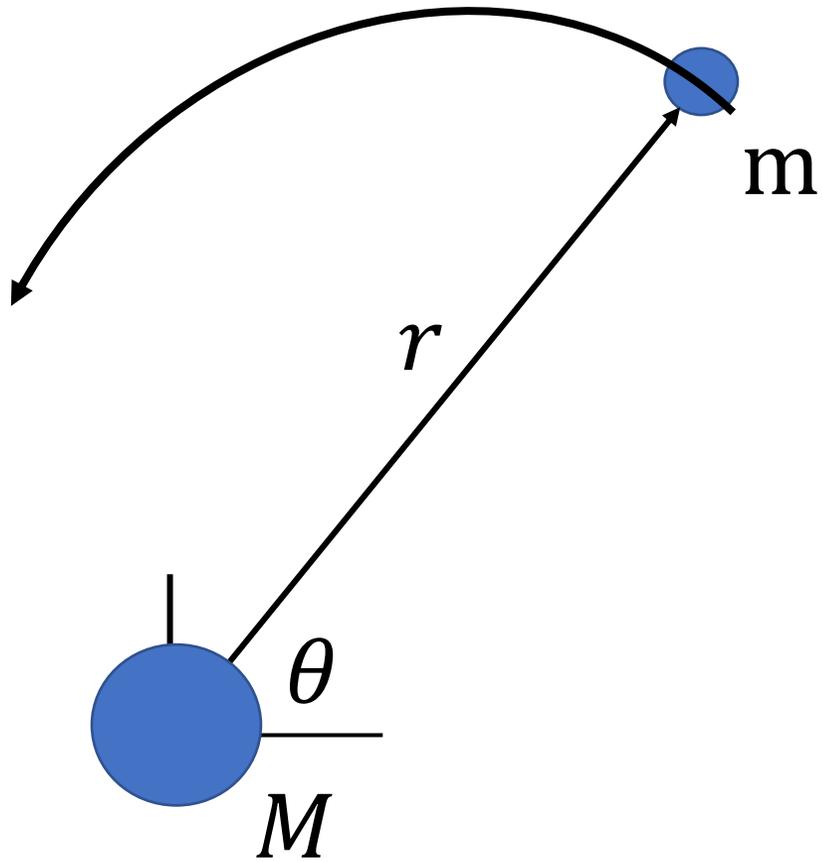


chain rule

$$\frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = r^2 \frac{d^2\theta}{dt^2} + 2r \frac{d\theta}{dt} \frac{d\theta}{dt}$$

multiply by r

$$\frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = 0$$

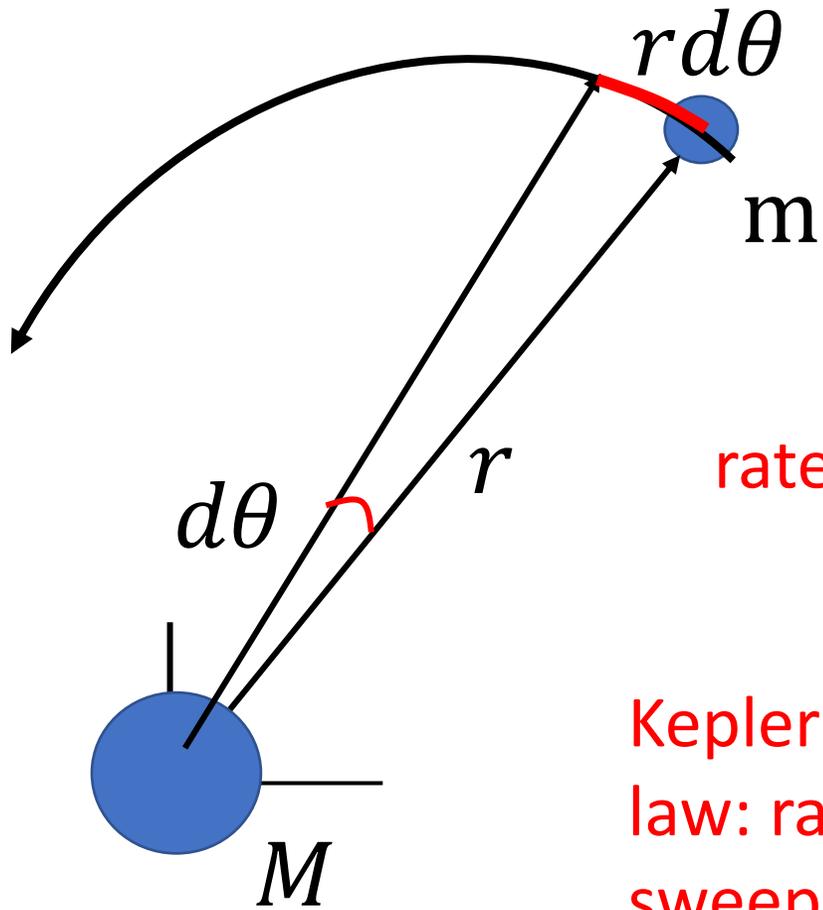


chain rule

$$\frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = r^2 \frac{d^2\theta}{dt^2} + 2r \frac{d\theta}{dt} \frac{dr}{dt}$$

so

$$r^2 \frac{d\theta}{dt} = \text{constant}$$



chain rule

$$\frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = r^2 \frac{d^2\theta}{dt^2} + 2r \frac{d\theta}{dt} \frac{dr}{dt}$$

rate of change of area = $r^2 \frac{d\theta}{dt}$

Kepler's second law: rate of sweeping out area is constant

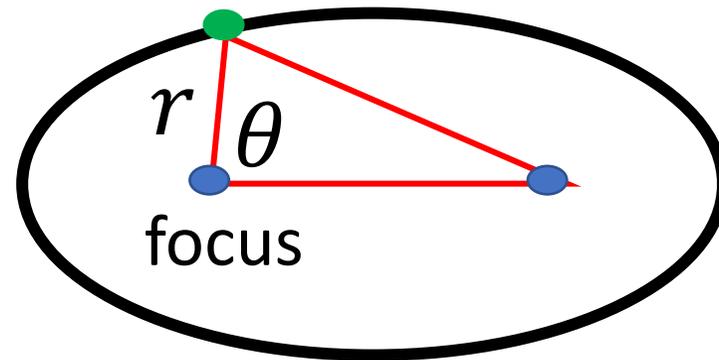
$$r^2 \frac{d\theta}{dt} = \text{constant}$$

further analysis shows that the shape of an orbit is an ellipse

formula for ellipse

$$r = \frac{p}{1 + \varepsilon \cos \theta}$$

p and ε constants



ellipse: length of red line
constant as green dot
moves around circumference

But the Earth is not a “point mass”



This issue stymied Newton for 20 years!

what Newton discovered:

The gravitational force outside a spherically-symmetric object is the same as for a point mass located at the object's center

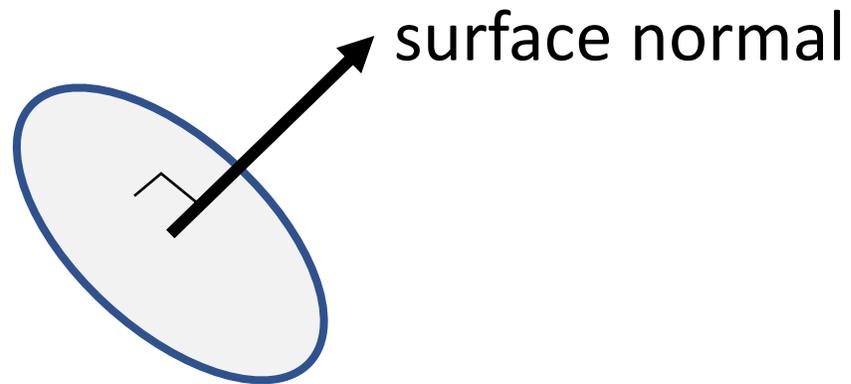
what Newton discovered:

The gravitational force outside a spherically-symmetric object is the same as for a point mass located at the object's center

I'll show you its true
but using modern concepts
that you are familiar with
(but Newton wasn't)

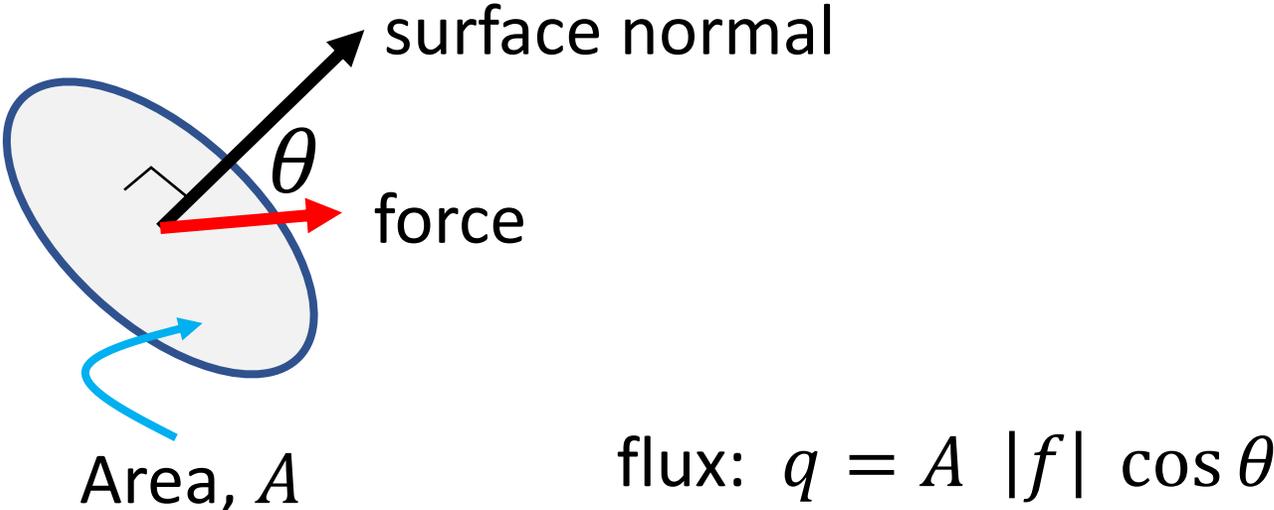
But first

I need to introduce the concept of
flux of force through a surface



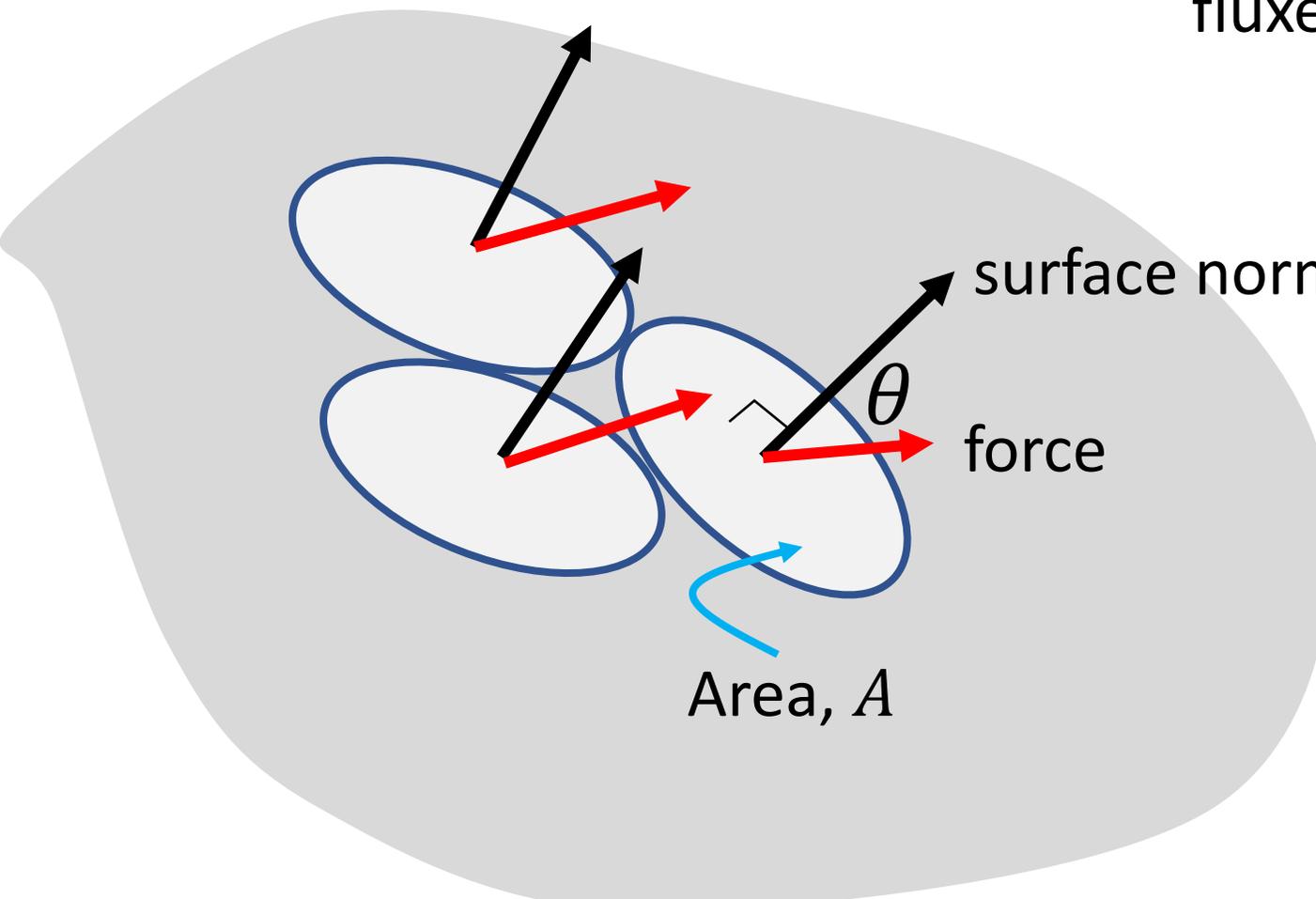
(this is the same as
water flux thru a surface
heat flux thru a surface)

for a small surface



for a big surface

just break into small
surfaces and add up the
fluxes of each



surface normal

force

Area, A

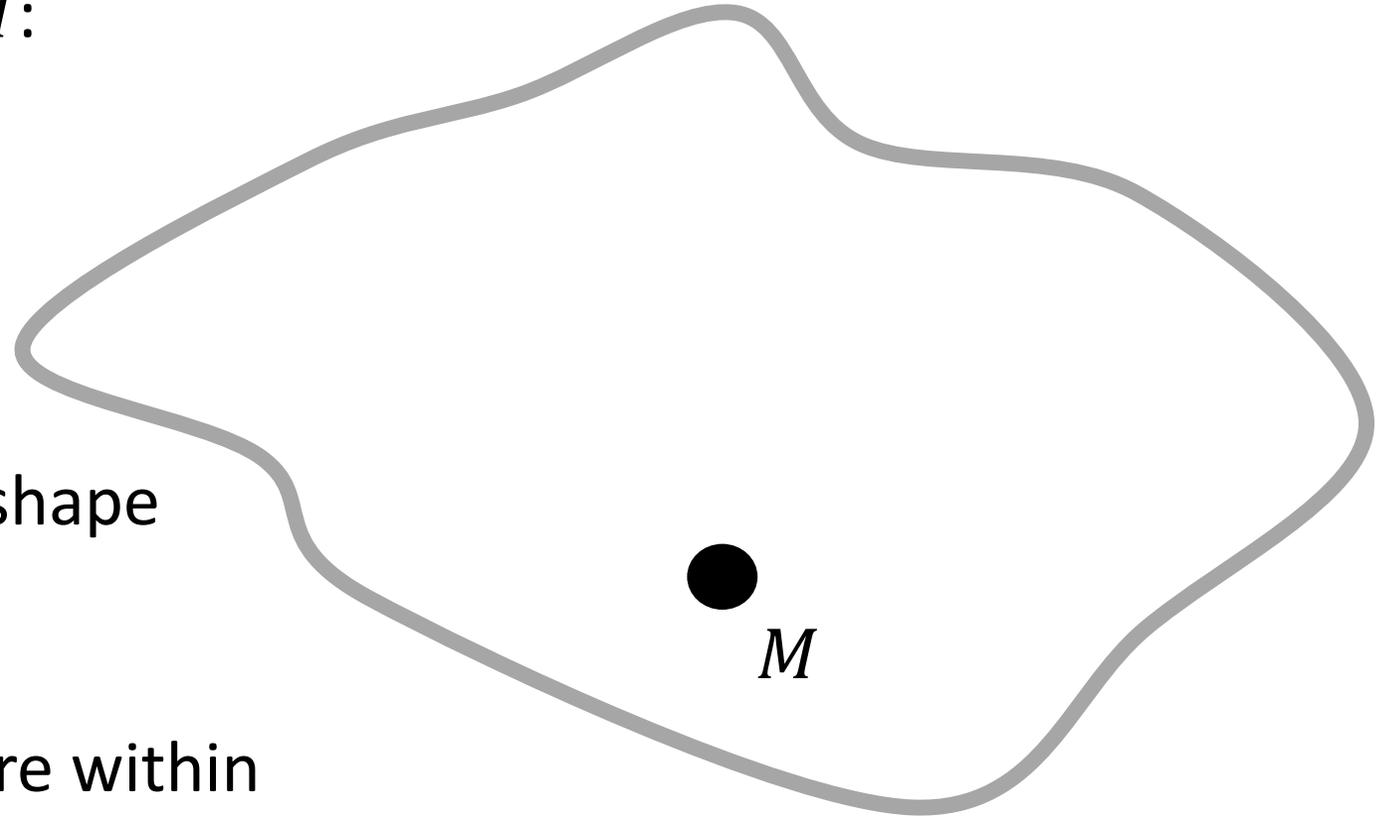
$$\text{flux: } q = \sum A |f| \cos \theta$$

Step 1:

The flux through a surface enclosing a point mass, M , is always $q = -4\pi\gamma M$:

irrespective of the shape
of the surface
and

irrespective of where within
the enclosed volume the point mass is located



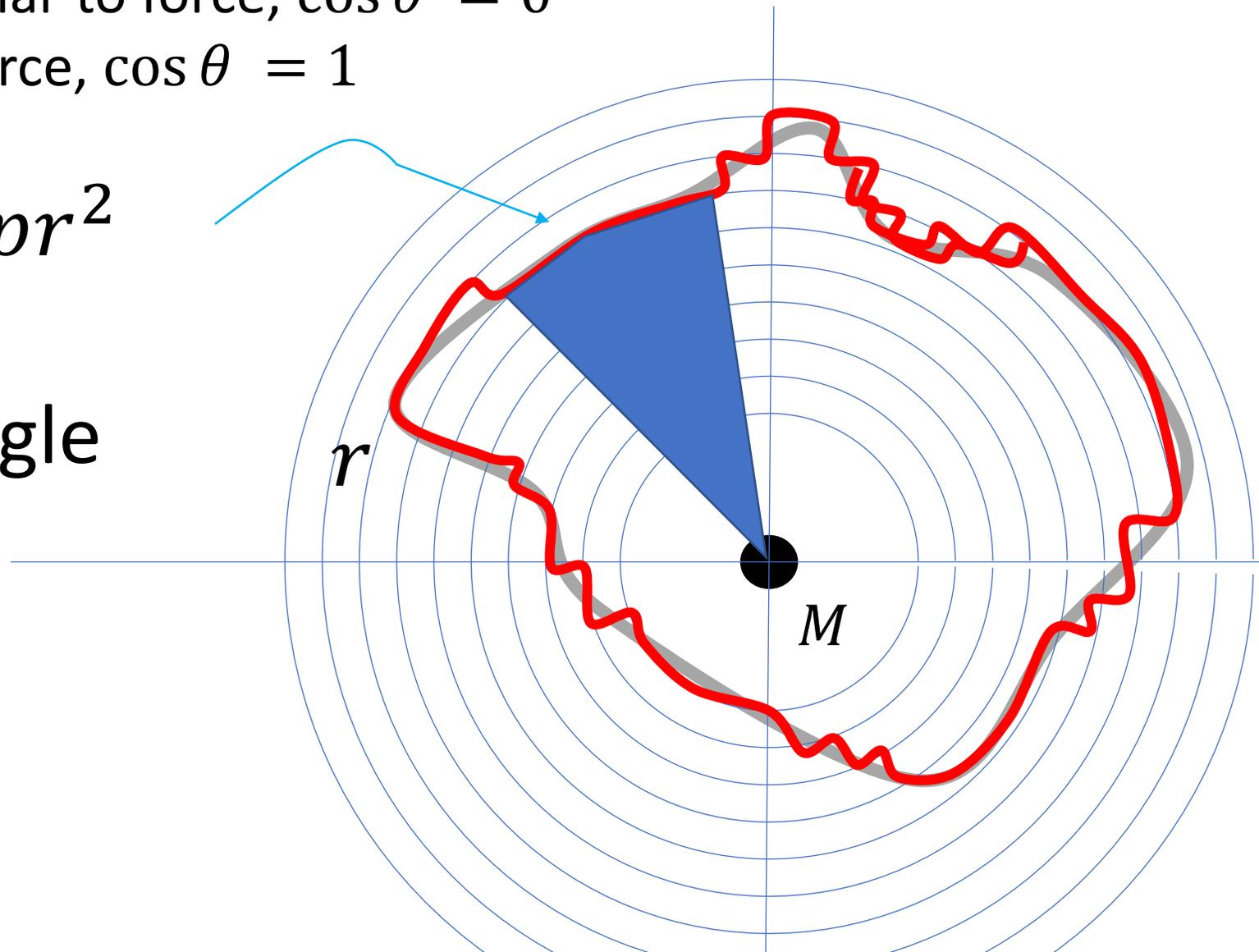
approximate the surface as sum of many “sides” and “caps”

“sides” normal perpendicular to force, $\cos \theta = 0$

“caps” normal parallel to force, $\cos \theta = 1$

area of a cap , $4\pi r^2 p$

p = fraction of
total spherical angle



approximate the surface as sum of many “sides” and “caps”

“sides” normal perpendicular to force, $\cos \theta = 0$

“caps” normal parallel to force, $\cos \theta = 1$

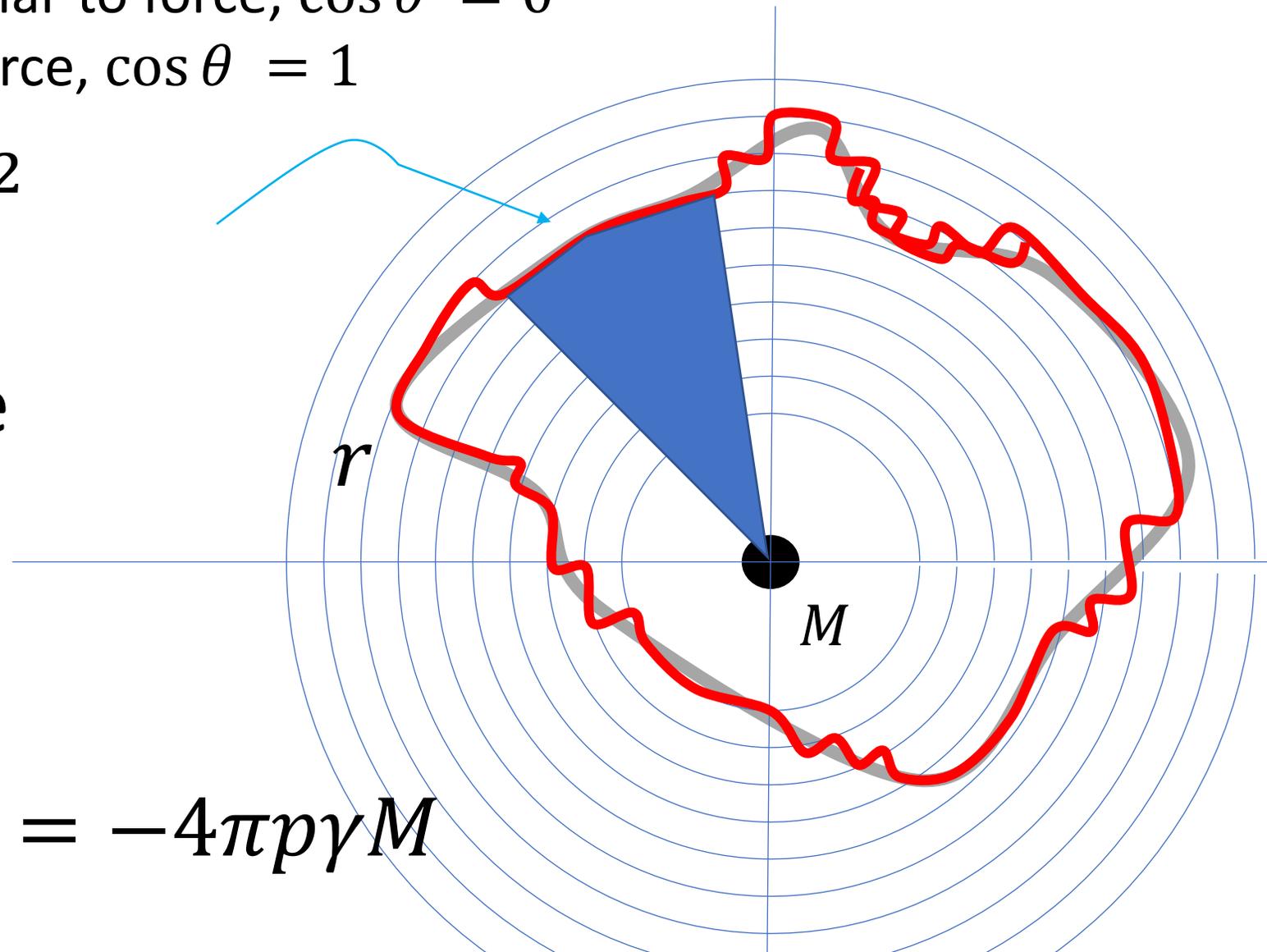
area of a cap , $4\pi pr^2$

p = fraction of
total spherical angle

force on cap , $-\frac{\gamma M}{r^2}$

flux thru cap

$$q = (4\pi pr^2) \left(-\frac{\gamma M}{r^2} \right) = -4\pi p\gamma M$$



approximate the surface as sum of many “sides” and “caps”

“sides” normal perpendicular to force, $\cos \theta = 0$

“caps” normal parallel to force, $\cos \theta = 1$

area of a cap , $4\pi p r^2$

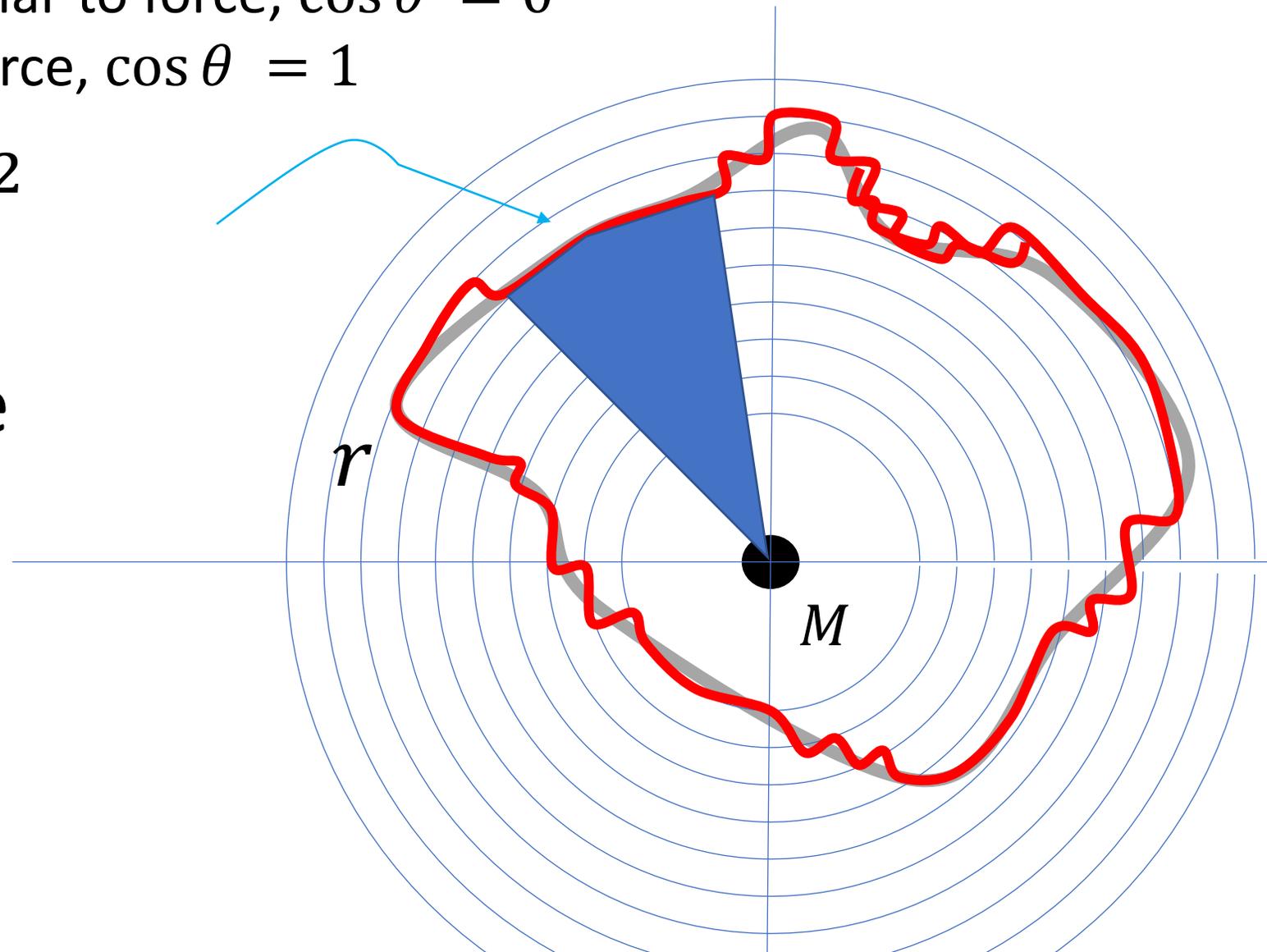
p = fraction of
total spherical angle

force on cap , $-\frac{\gamma M}{r^2}$

flux thru all caps

$$q = -4\pi\gamma M$$

size p 's sum to 1



Step 2:

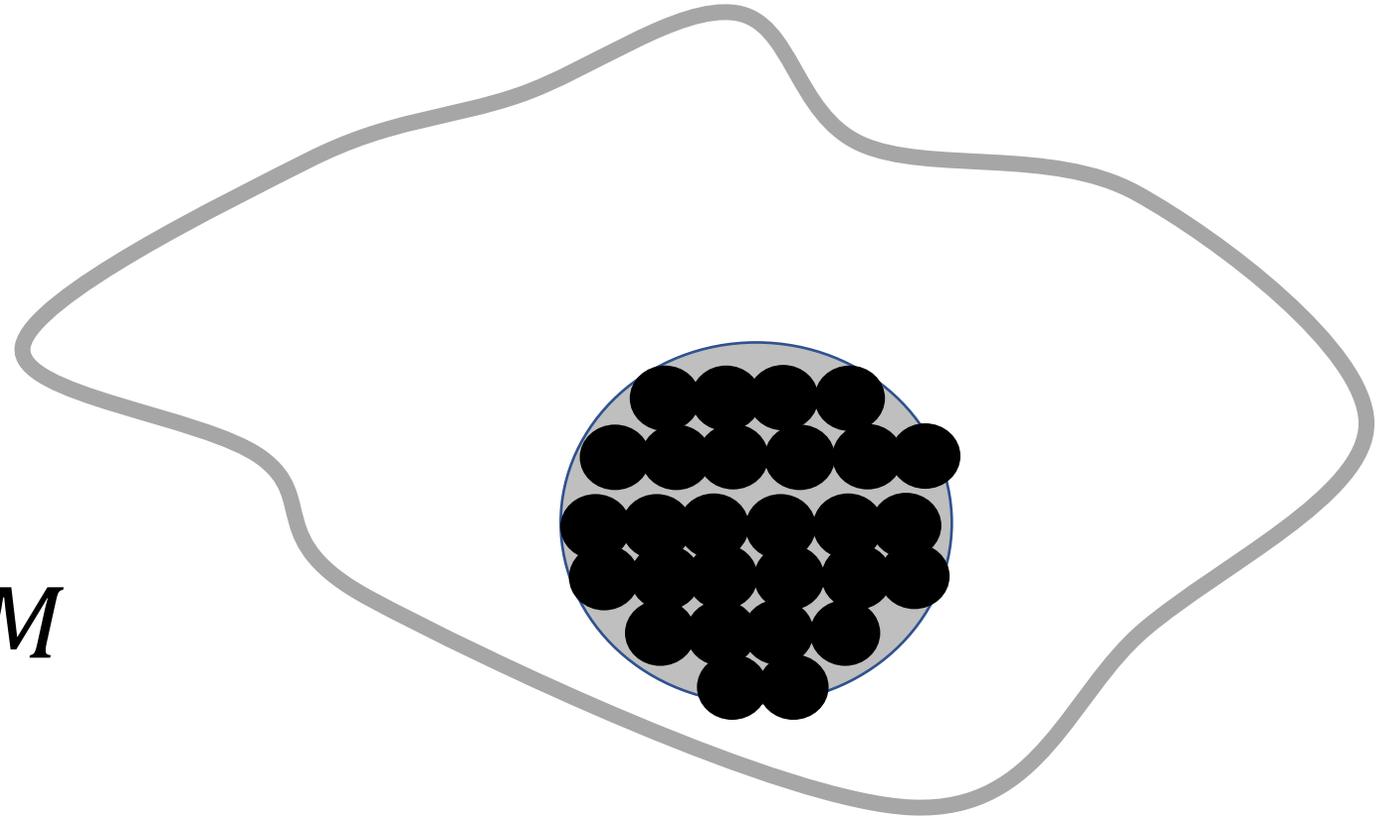
construct sphere by assembling many point masses with total mass, M

flux is still

$$q = -4\pi\gamma M$$

since fluxes sum
and

flux from any one point mass is independent of its position



Step 3:

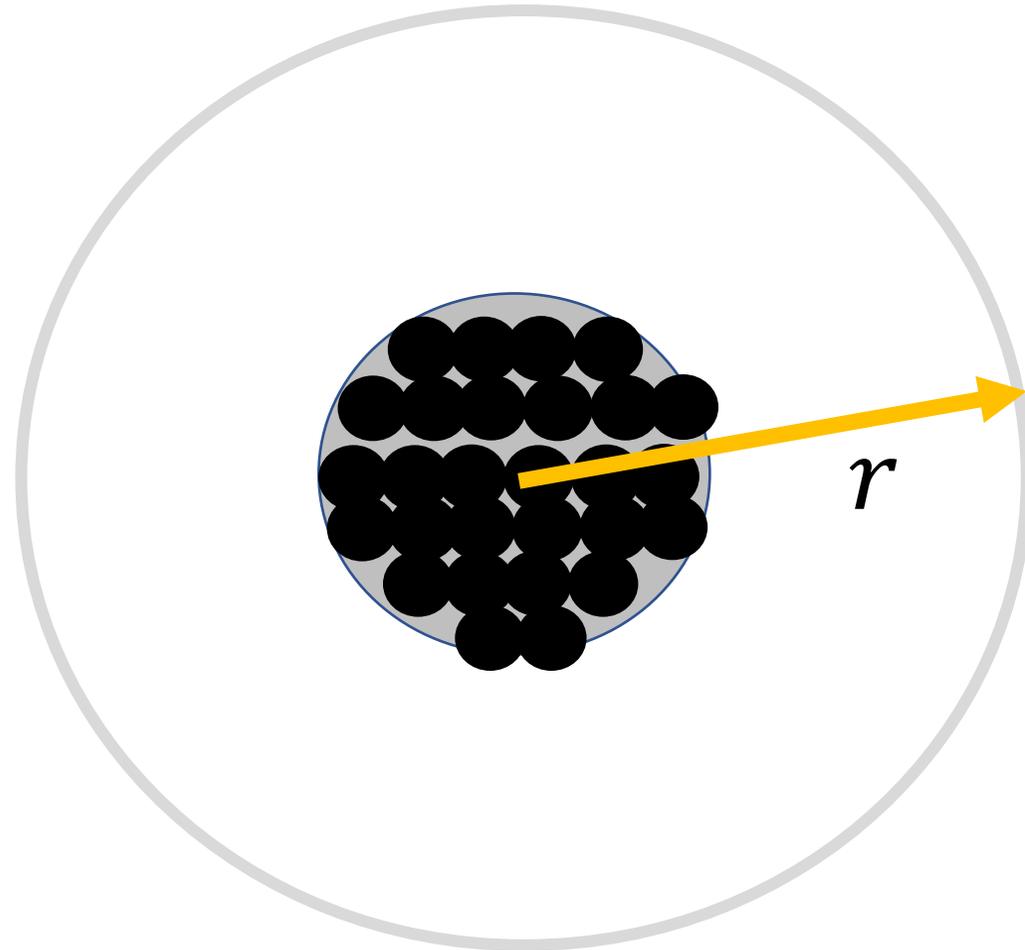
now make the surface a sphere of radius r centered at the center of the mass, so the problem is spherically symmetric

flux is still

$$q = -4\pi\gamma M$$

since flux is independent of shape of surface

furthermore, this shape is all “cap”, so $\cos \theta = 1$



Step 4:

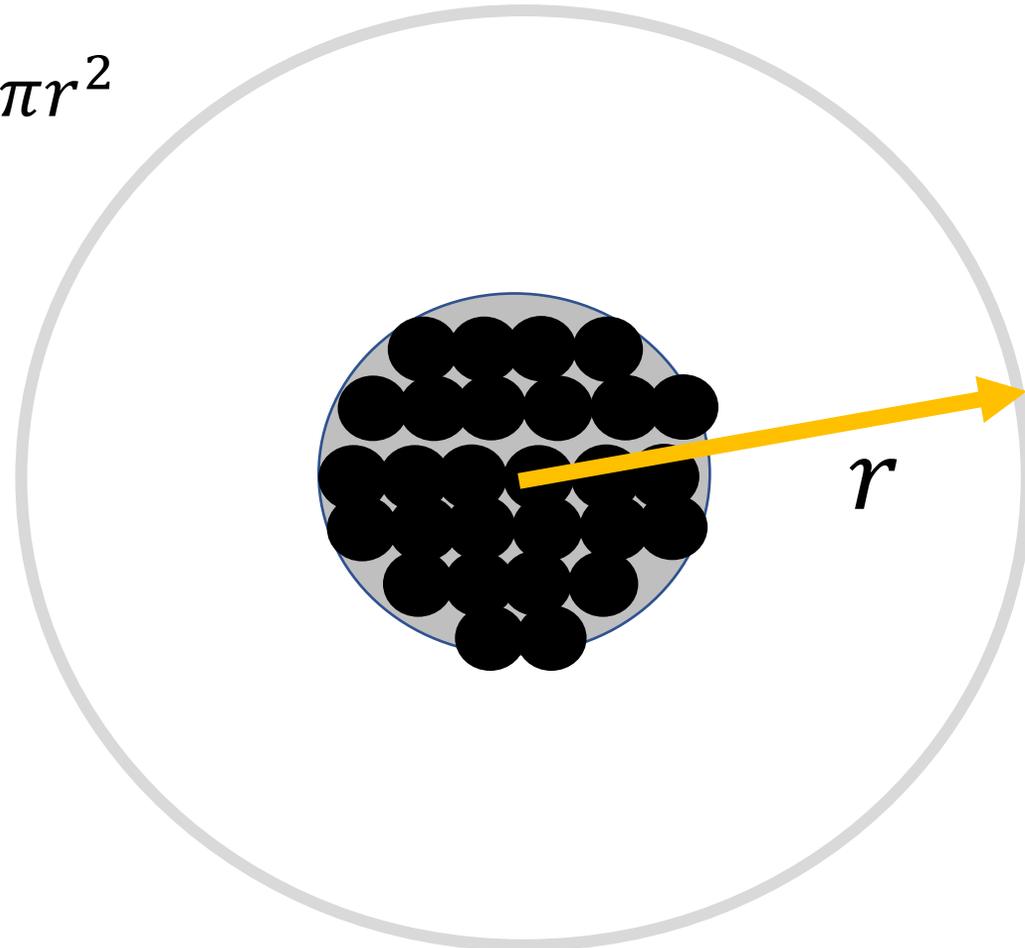
As the problem is spherically symmetric, the flux must be constant over the surface of the sphere

The sphere has area, $A = 4\pi r^2$

so the flux per area is

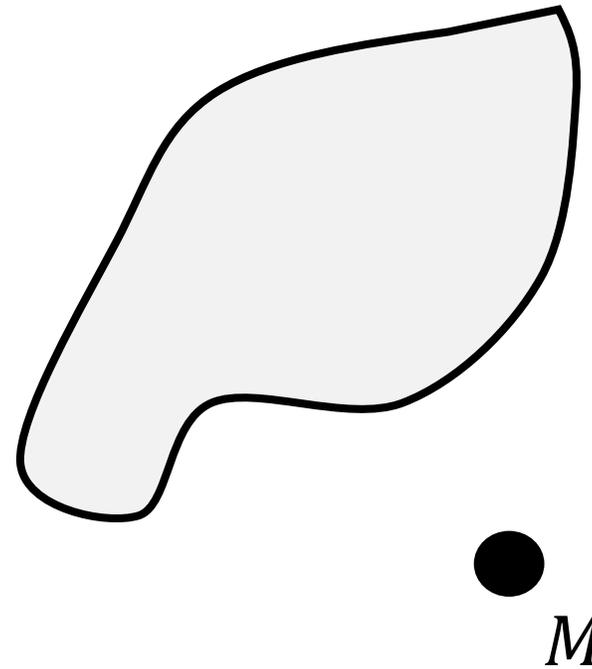
$$q/A = -\frac{\gamma M}{r^2}$$

which is equal to the force and has the form of the force from a point mass



An interesting tidbit

The flux through a surface excluding a point mass, M , is always $q = 0$:



irrespective of the shape
of the surface
and

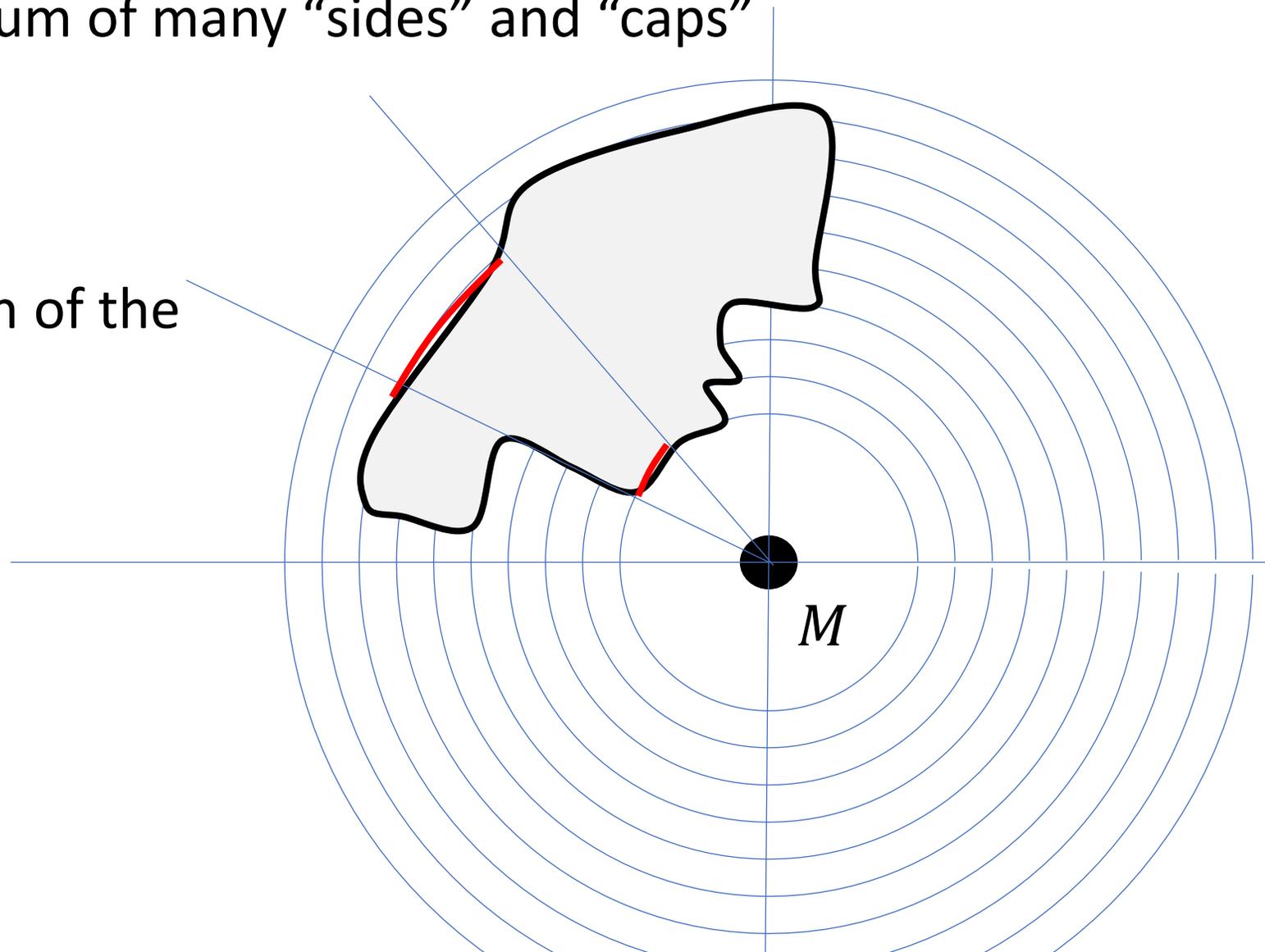
irrespective of where outside
the excluding volume the point mass is located

handled the same way

approximate the surface as sum of many "sides" and "caps"

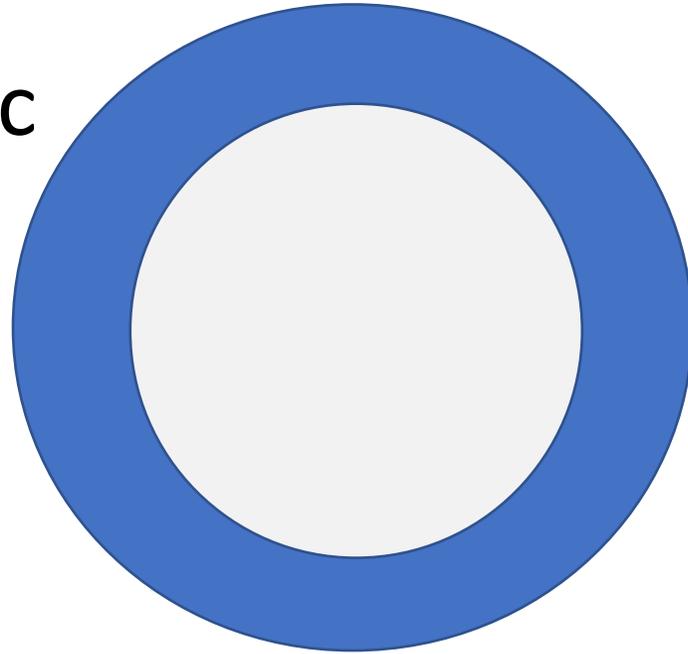
but now that caps are paired

you find that the contribution of the two caps cancel



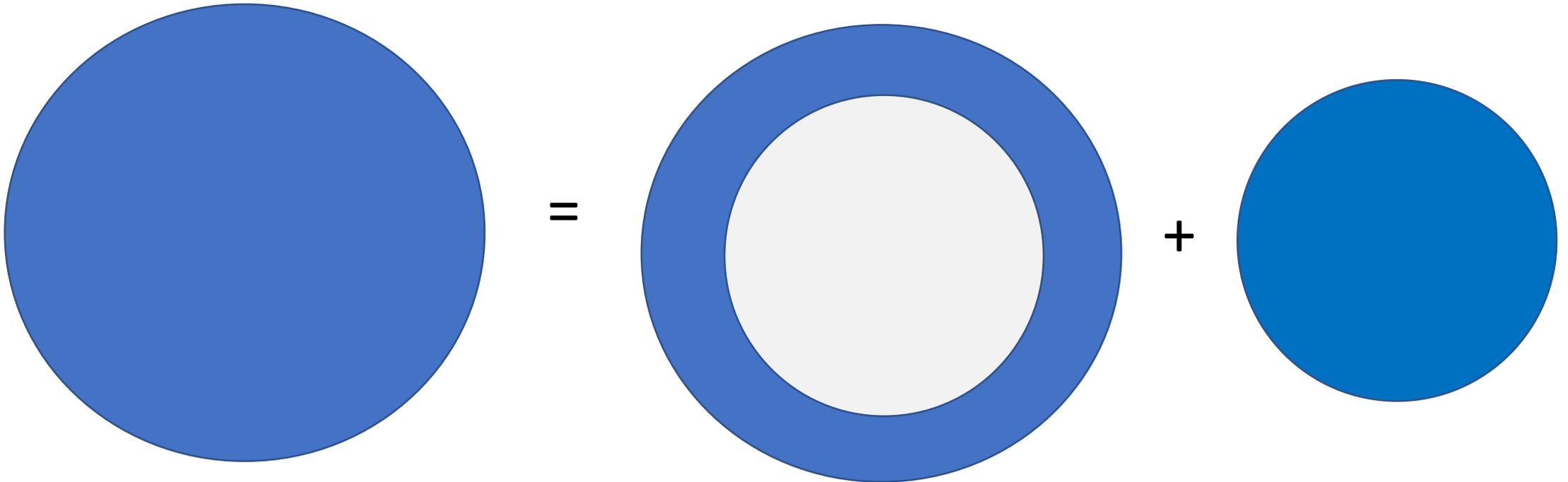
You can use this to show

force inside a cavity
inside a spherically symmetric
object is zero

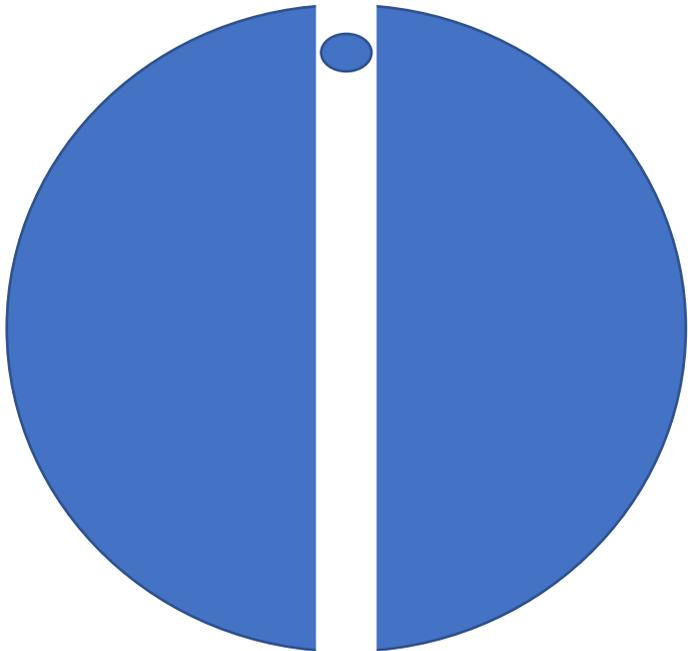


And

force at depth in a sphere depends only on mass below you



throw rock of mass m into small hole drilled thru earth



what do you think will happen?

throw rock of mass m into small hole drilled thru earth

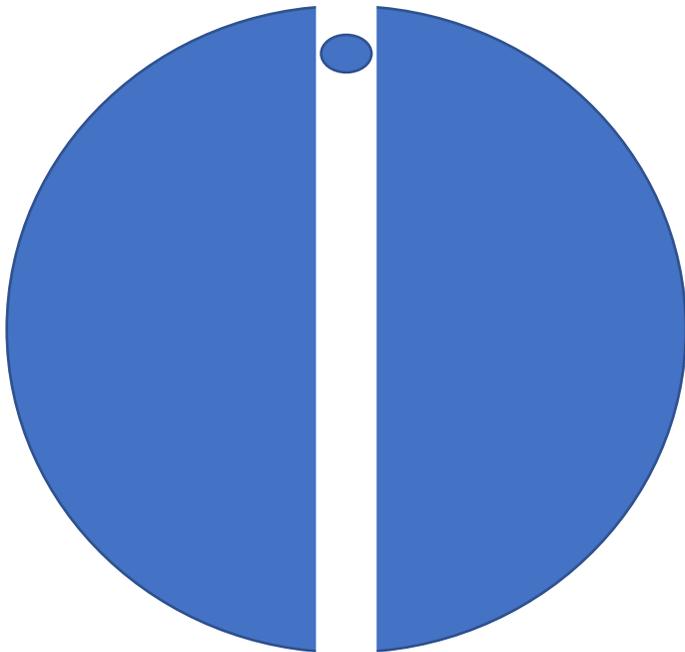
force law

$$f_r = m \left(\frac{d^2 r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 \right)$$

Newton's Law

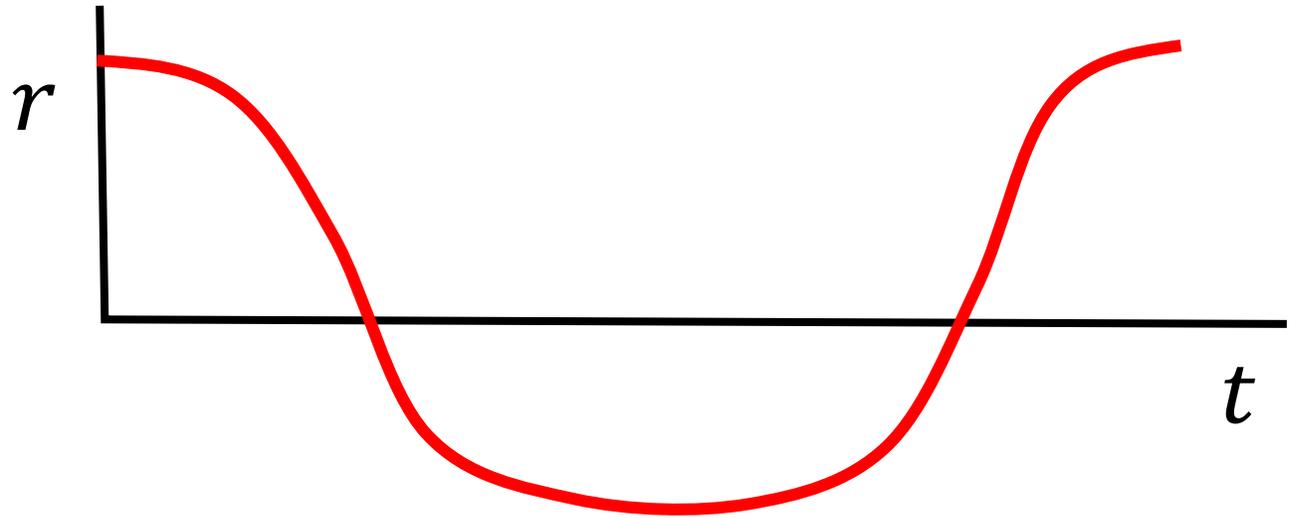
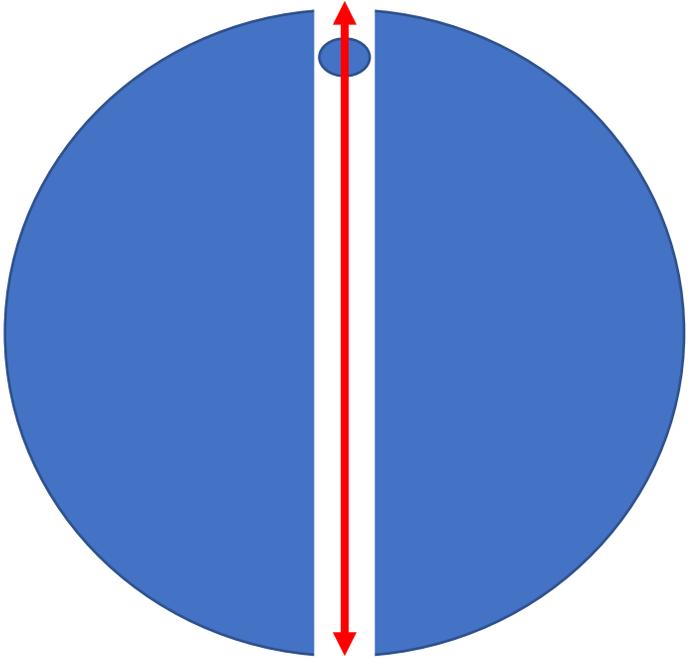
$$f_r = -\frac{\gamma m M(r)}{r^2} \quad \text{with } M(r) \approx \frac{4}{3} \pi r^3 \rho$$

$$f_r = -\frac{4}{3} \gamma m \pi r \rho = -C m r$$



$$\frac{d^2 r}{dt^2} = -Cr$$

$$r = R_0 \cos(\sqrt{C}t)$$



measuring gravity

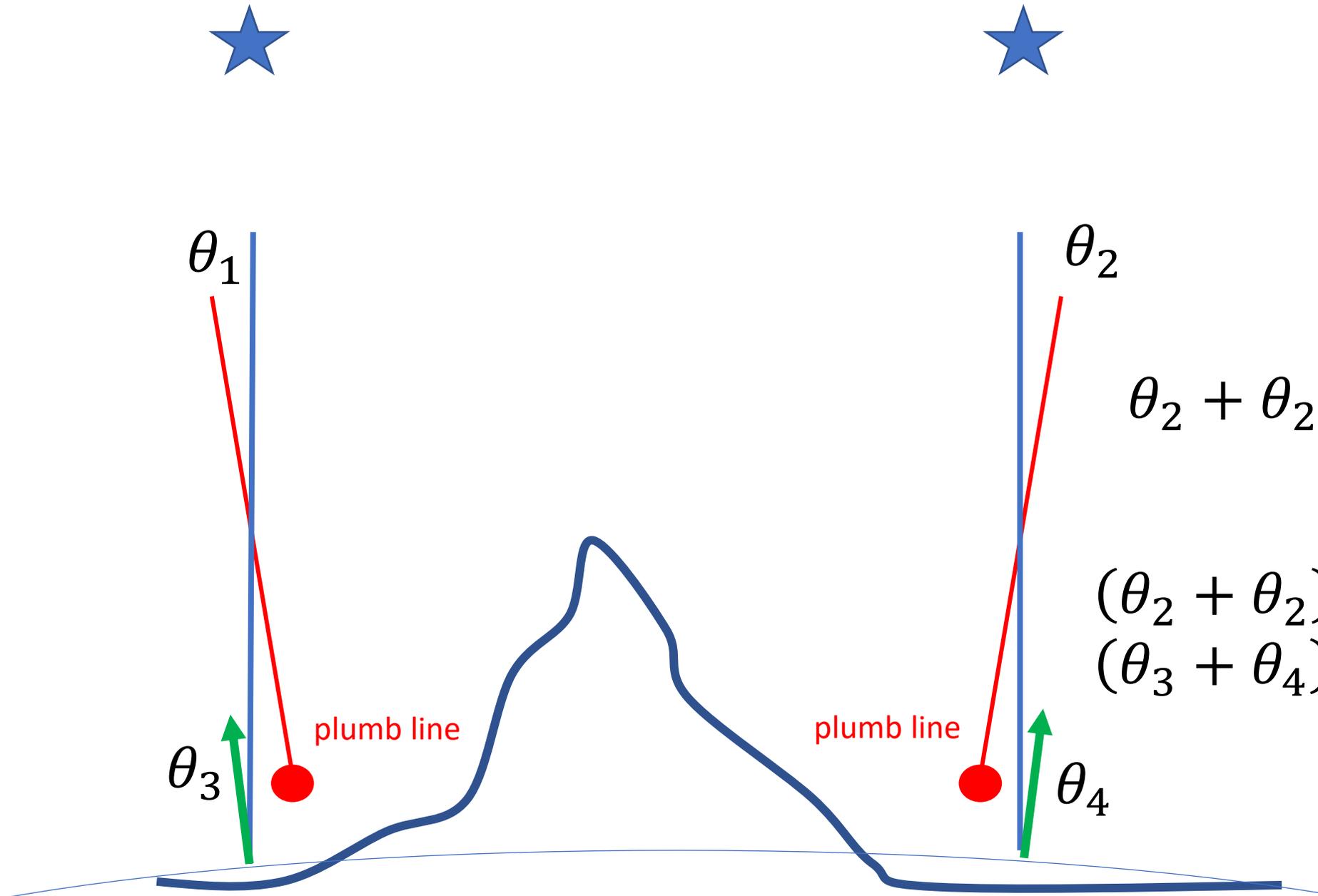
does gravity point straight down?

and if it doesn't

how would you tell?

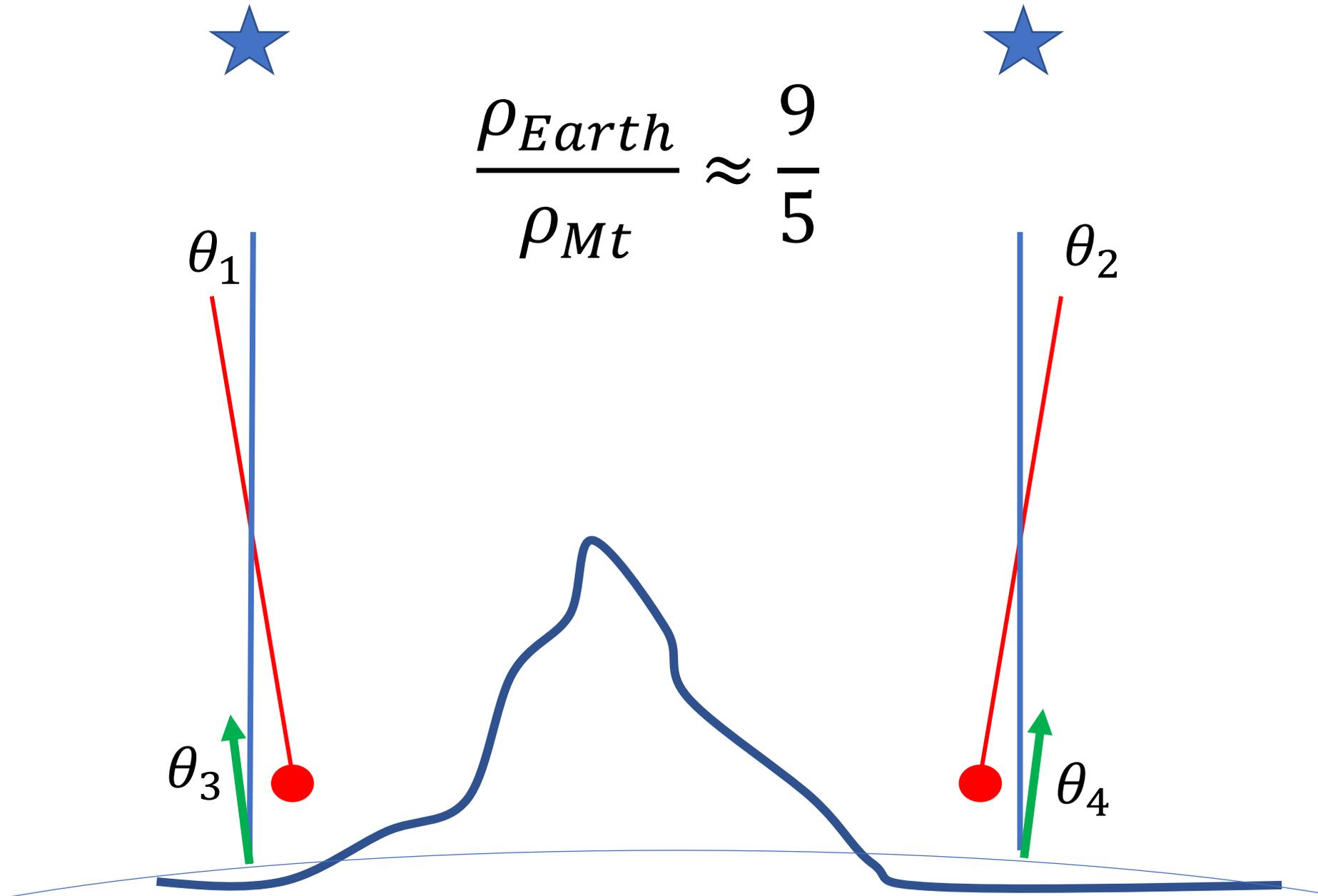


Mt Schiehallion Experiment, 1774



$$\theta_2 + \theta_2 = 54.6''$$

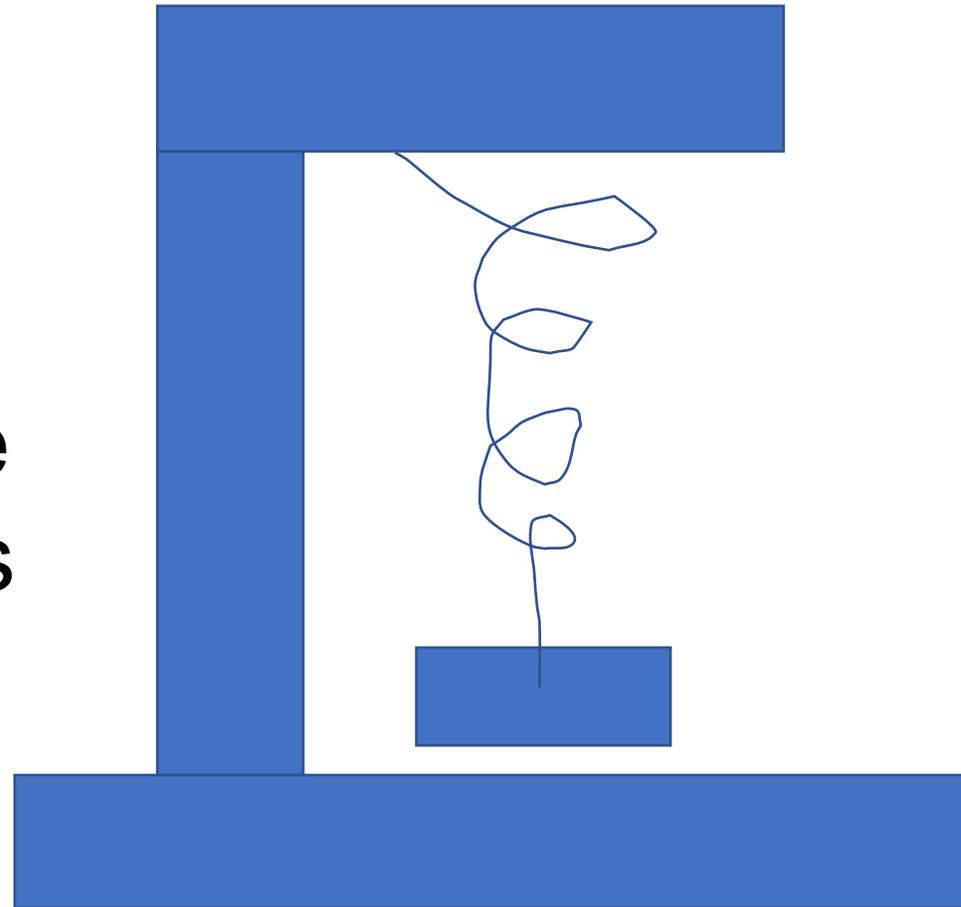
$$(\theta_2 + \theta_2) - (\theta_3 + \theta_4) = 11.6''$$



How do you measure the strength of gravity?

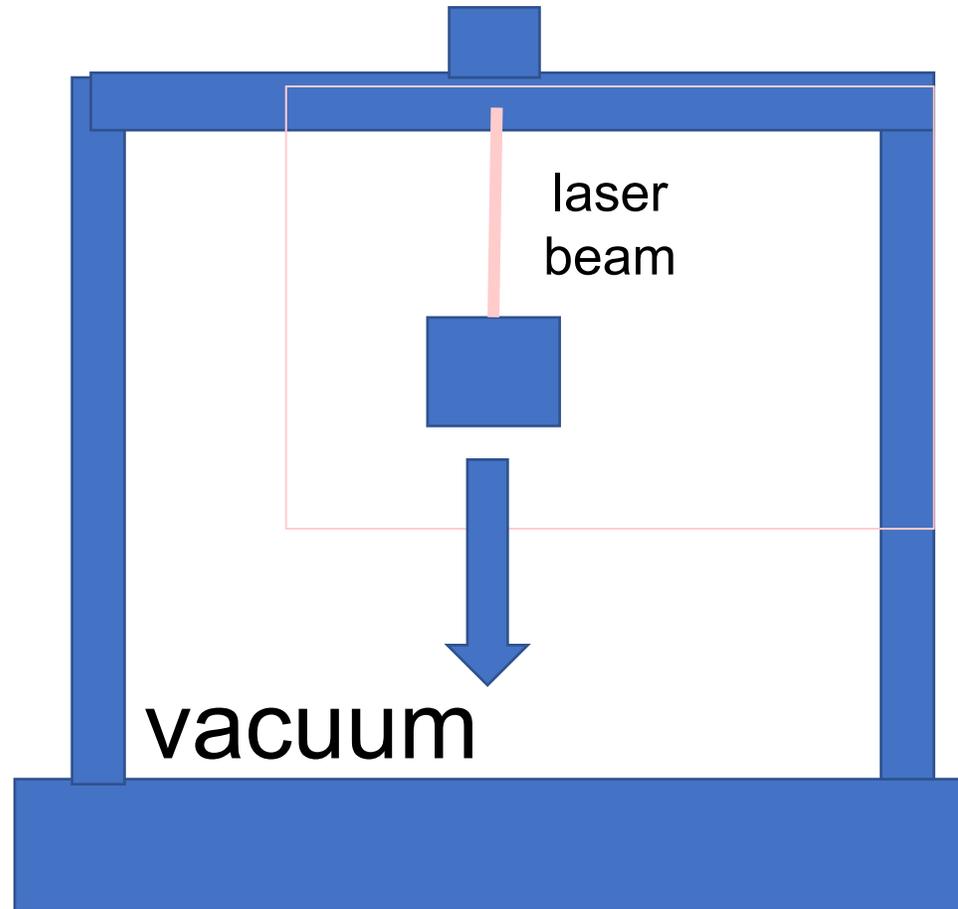
deflection of a mass weighing down a spring

best at relative
measurements



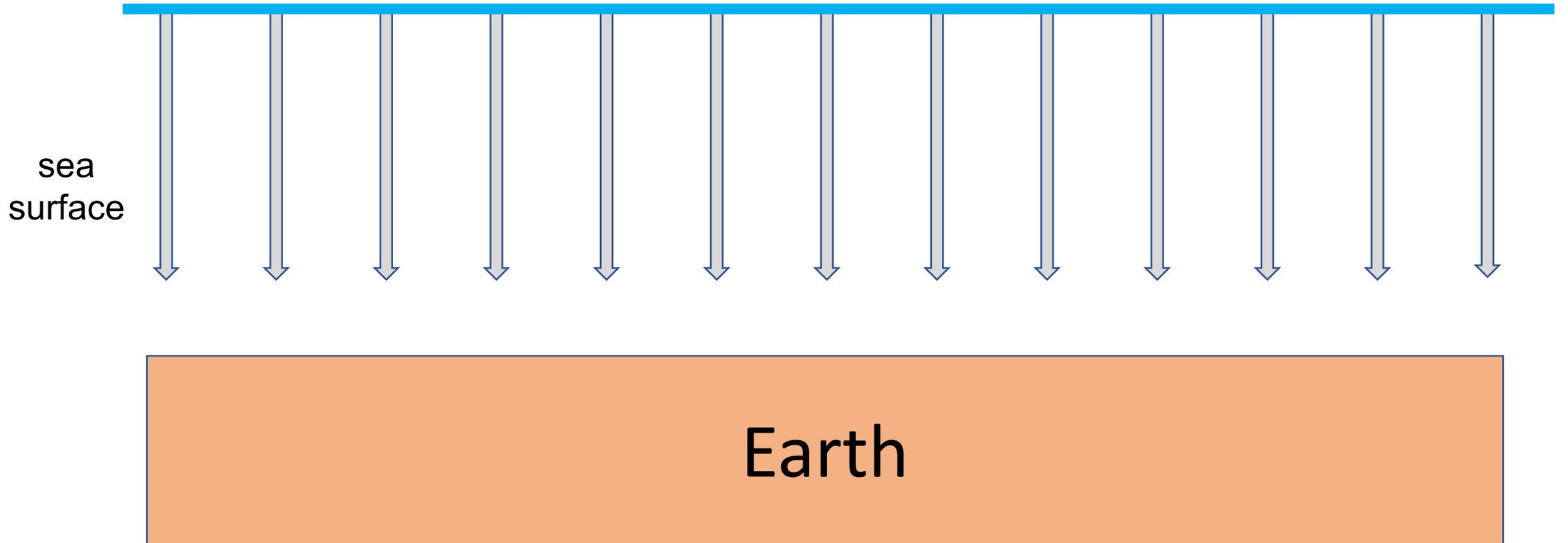
measure
small
changes
in height

drop small mass in vacuum and measure its
acceleration

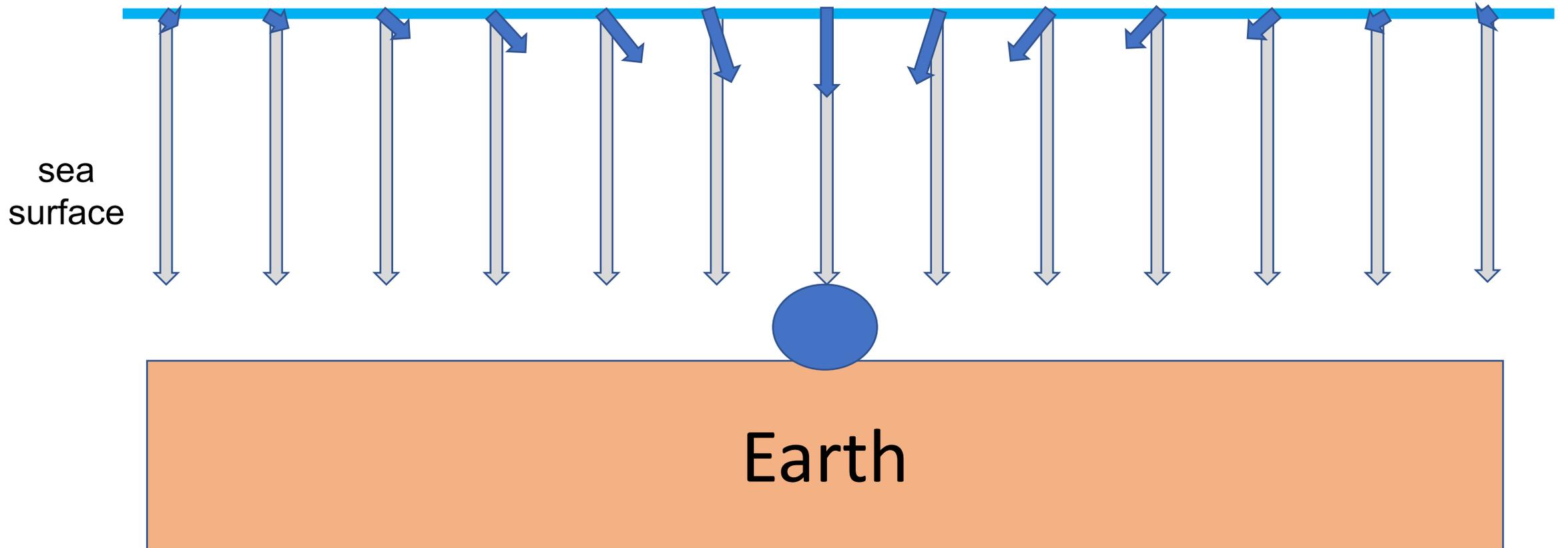


measure
acceleration
optically
with
laser
interferometry

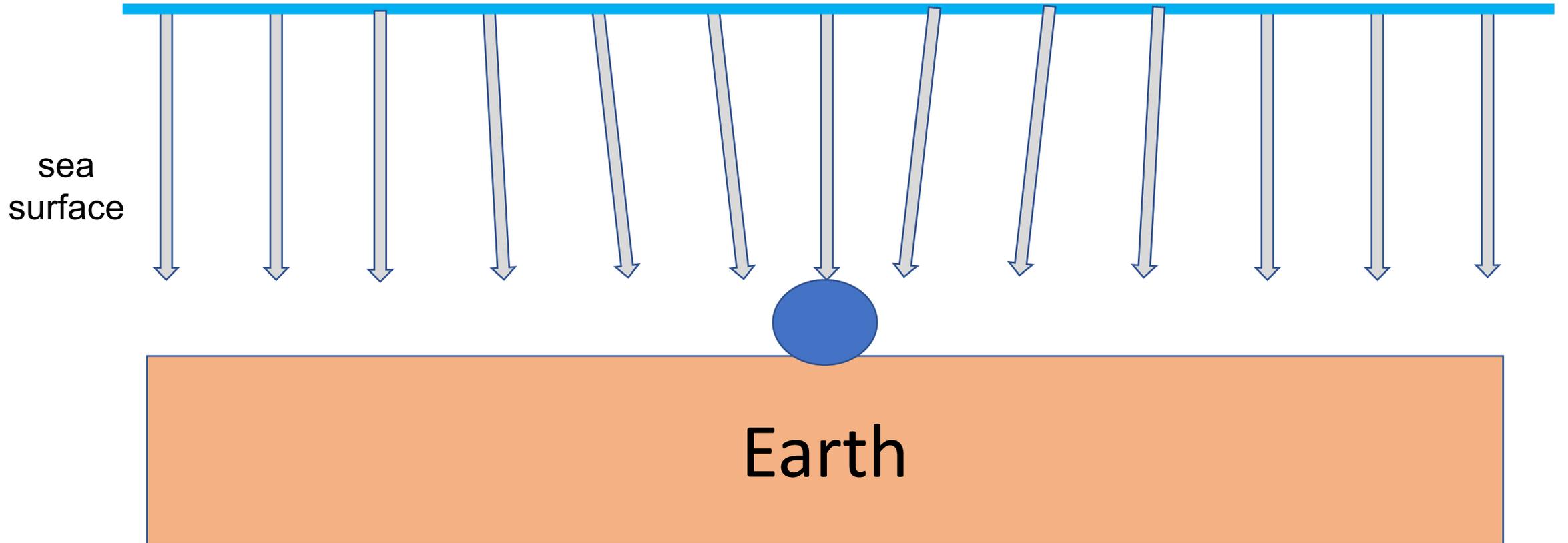
A fluid adjusts so its surface is everywhere parallel to the local gravitational field



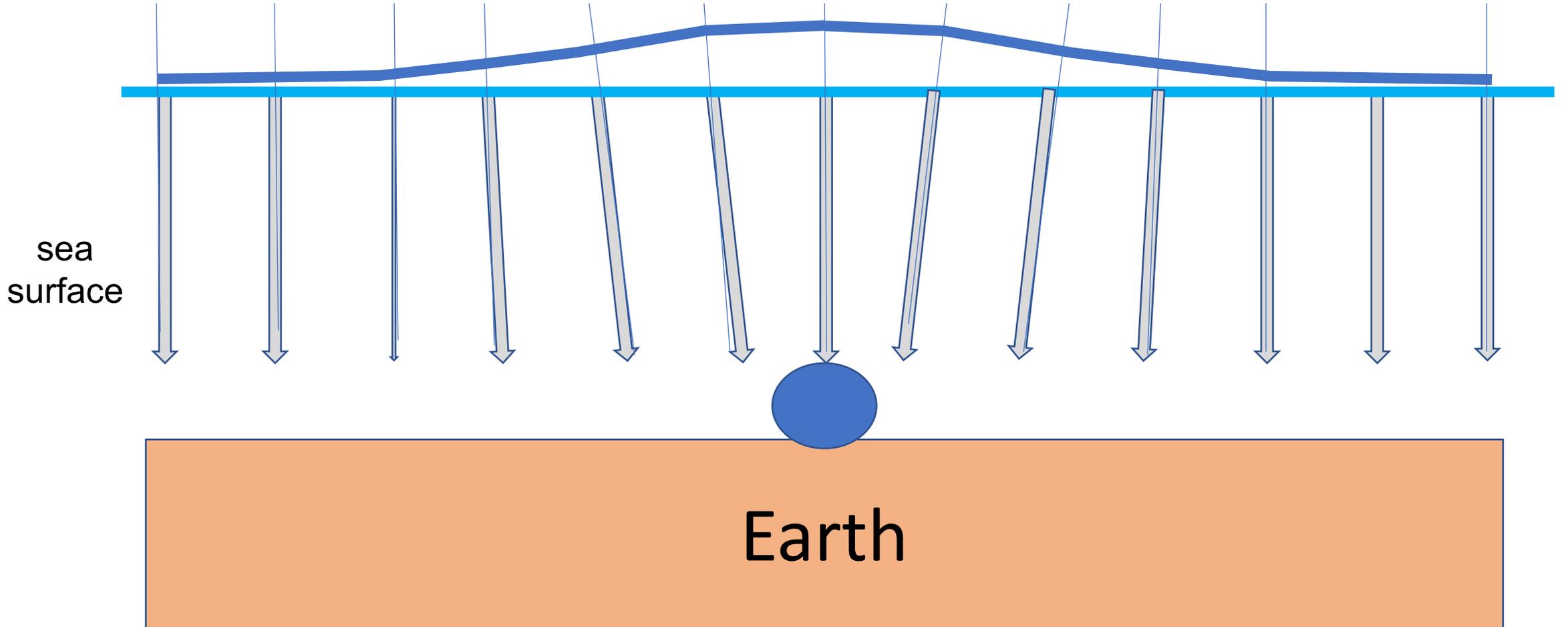
A fluid adjusts so its surface is everywhere parallel to the local gravitational field



Which way does the water go?



bulges up





What will happen to
sea level near
Greenland when
the ice melts?

